

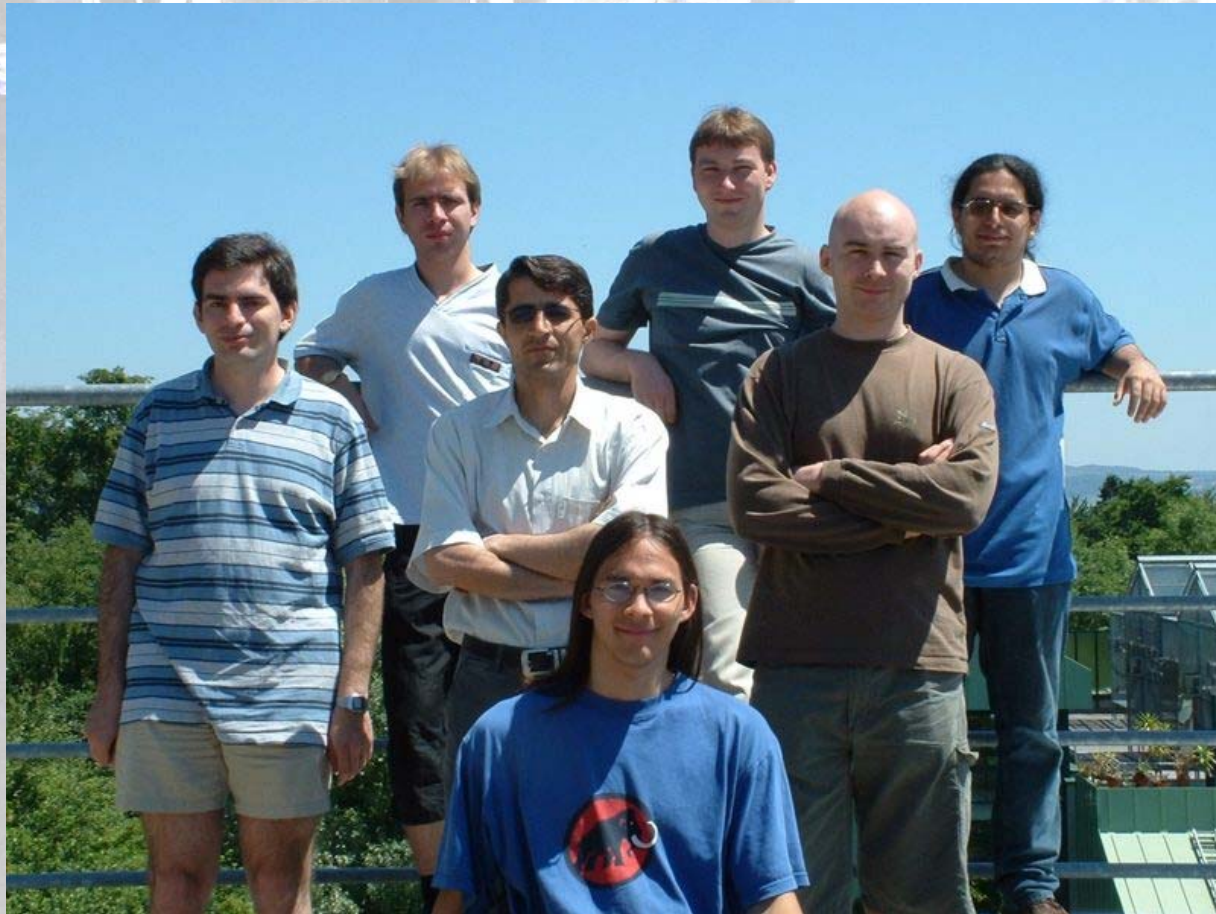


***Quantum information storage in  
atomic media***

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**Alex Lvovsky's Quantum Technology Group**

# Overview

## Introduction

- What is Electromagnetically induced transparency?
- Setting up the EIT experiments
- Experimental EIT
- Measuring the EIT width
- Ground state decoherence
- Storage of light
- Single photon source

## Conclusions

# Introduction

- Light is a good carrier of quantum information.
- Every computer needs memory.
- Photons don't like to stay put.
- Atoms are very good to store information.
- Atoms interact strongly with light.

**Our task is to develop an interface that would allow transfer of quantum information between optical and atomic media.**

# Introductory terms

## Group velocity

$$v_g(\omega) = \frac{c}{n(\omega) + \omega \frac{dn}{d\omega}}$$

This term can allow us to reduce the group velocity

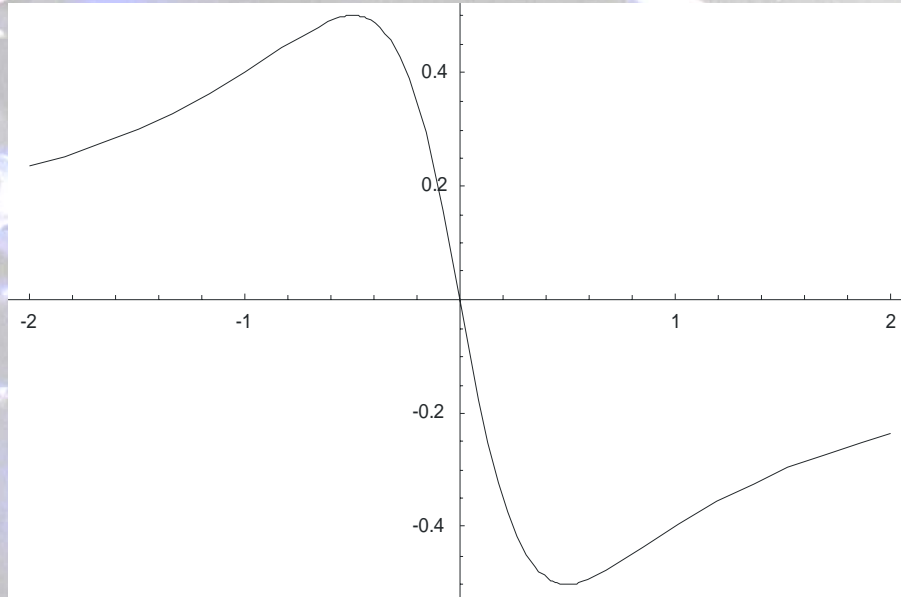
## Dispersion relations

$$n = n_r + i \cdot \frac{\alpha}{2\frac{\omega}{c}}$$

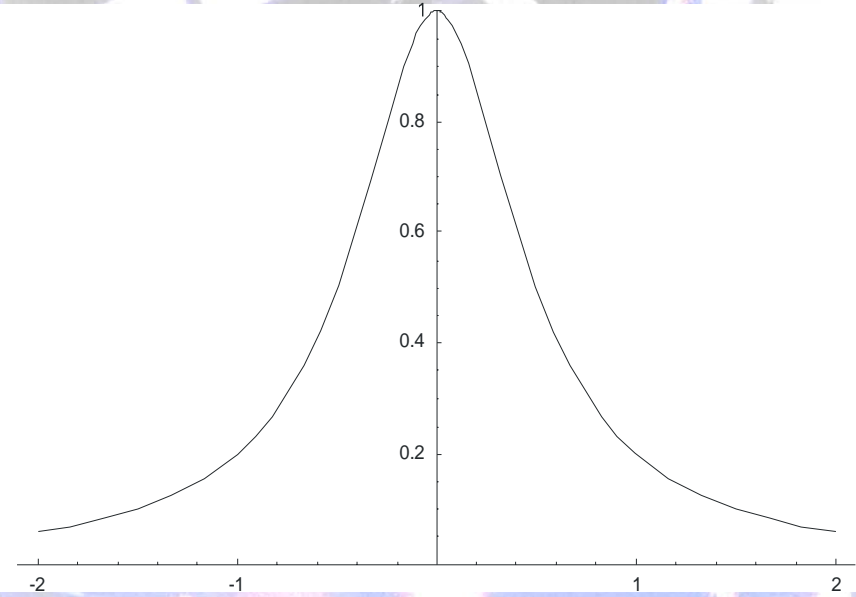
Dispersion	$n_r = \text{Re}(n)$
Absorption	$\alpha = \frac{2\omega}{c} \text{Im}(n)$

# Behavior close to an atomic resonance

Dispersion



Absorption



**The problem: In normal media a change in the group velocity is always associated to big absorption**

# Introducing EIT

The solution: electromagnetically-induced transparency (EIT).

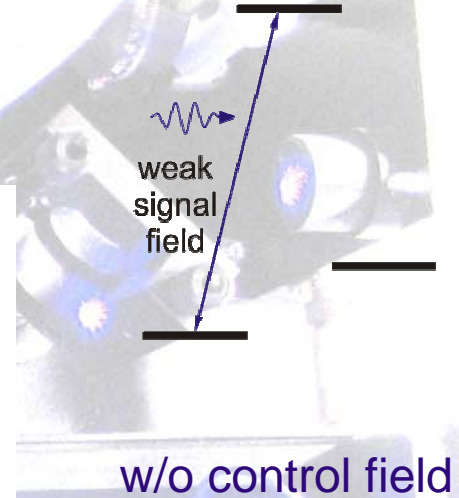
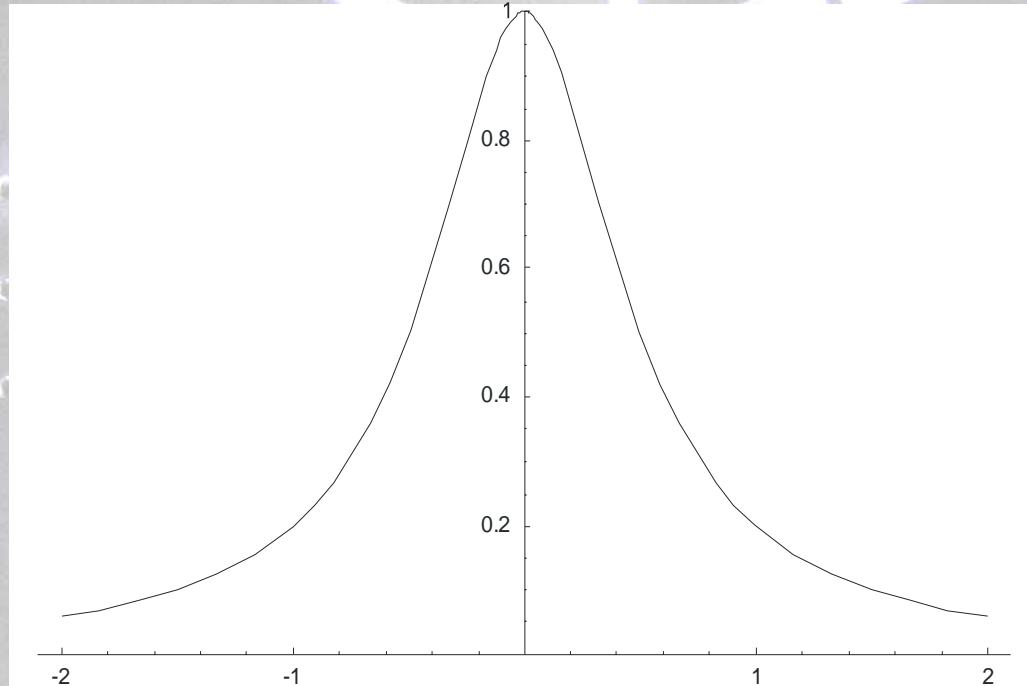
Lambda system

weak  
signal  
field

strong  
control  
field

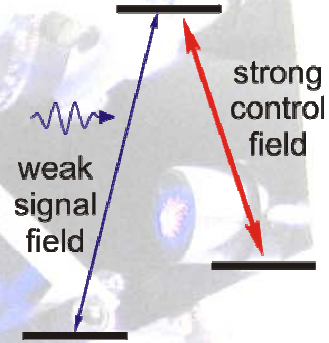
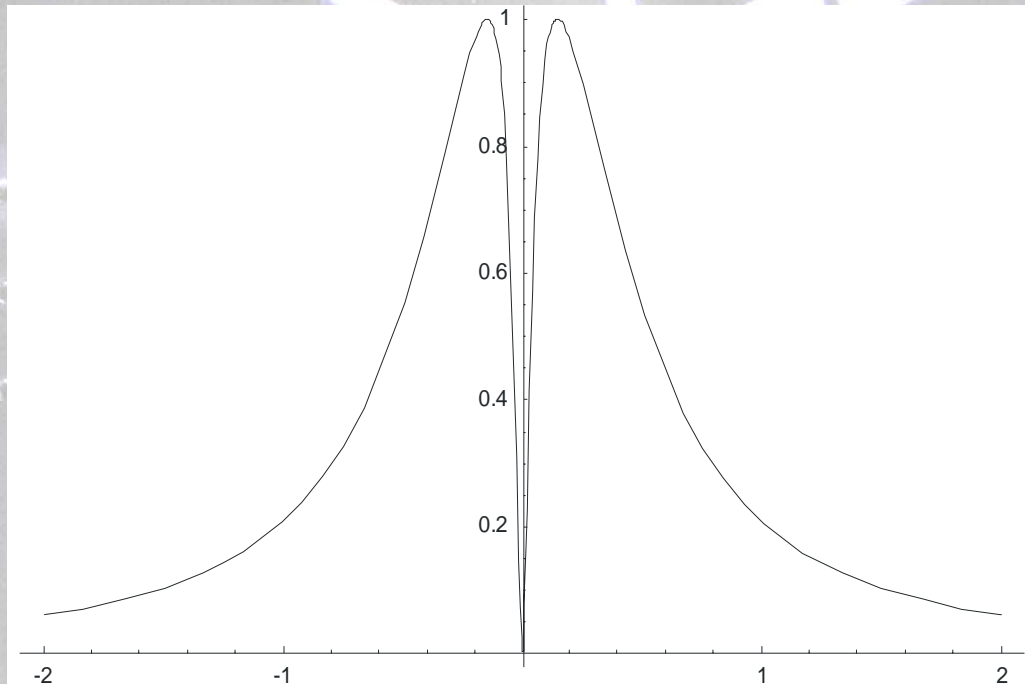
Harris S. Phys. Today 50 (7) 36-42 (1997)

# Absorption of the signal field





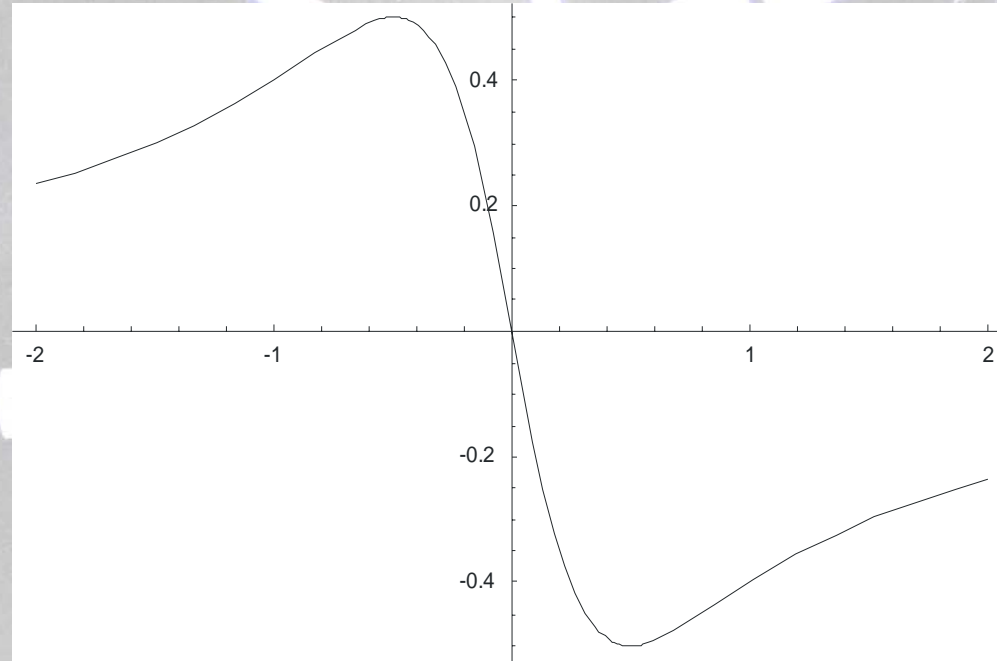
# Absorption of the signal field



with control field

- Light propagates through an otherwise opaque medium.

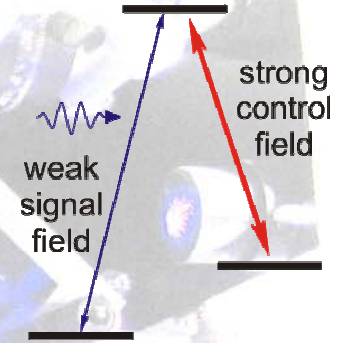
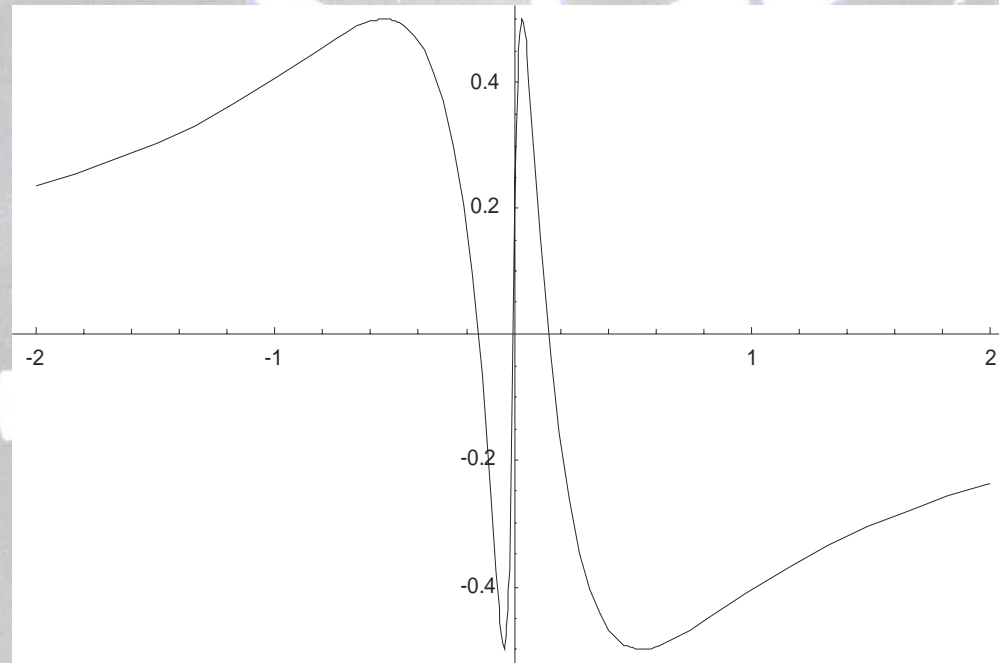
# Dispersion of the signal field



weak  
signal  
field

w/o control field

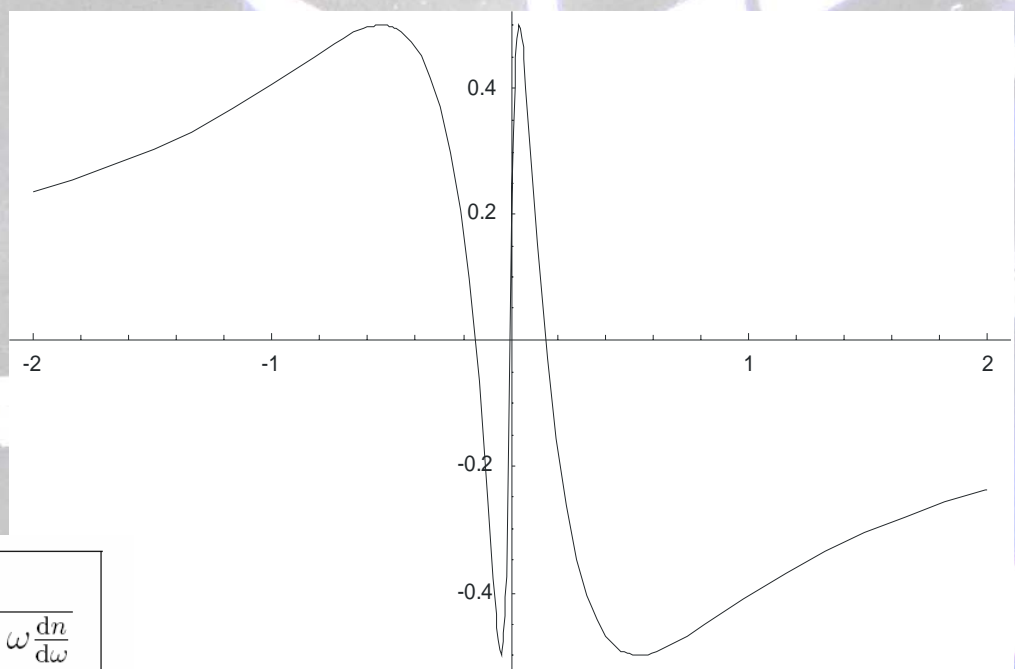
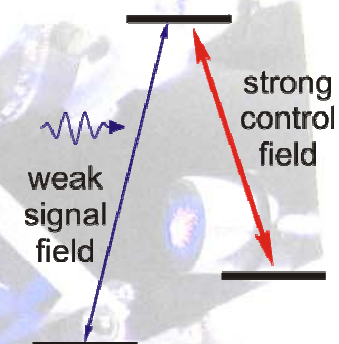
# Dispersion of the signal field



with control field

- We can enormously reduce the group velocity
- The slope is proportional to the strength of the control field

# Dispersion of the signal field



with control field

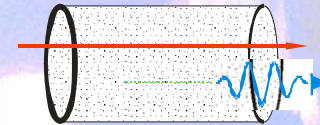
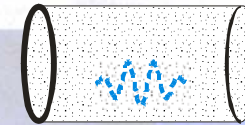
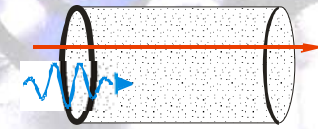
**Remember:**

$$v_g(\omega) = \frac{c}{n(\omega) + \omega \frac{dn}{d\omega}}$$

- We can enormously reduce the group velocity
- The slope is proportional to the strength of the control field

# How can we use EIT to store information?

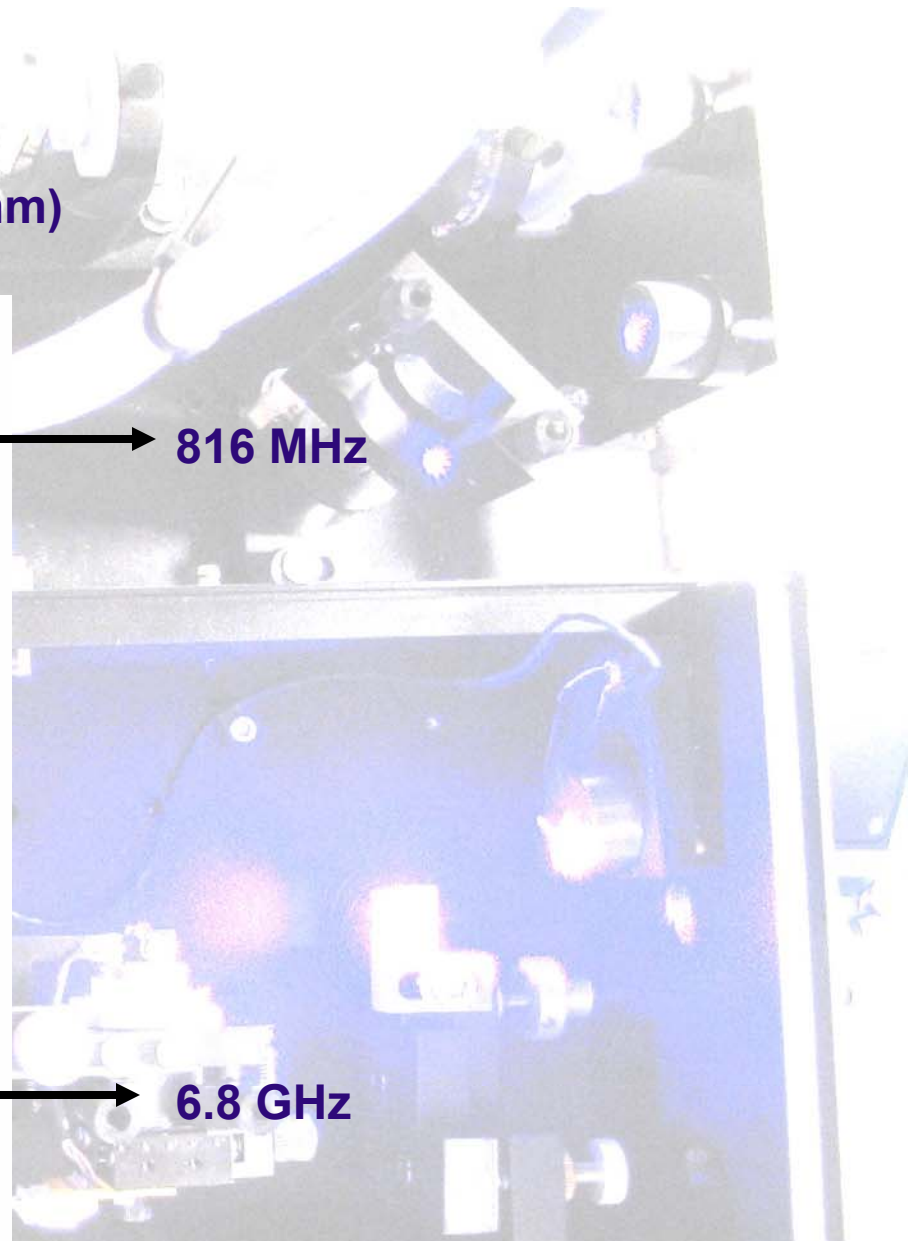
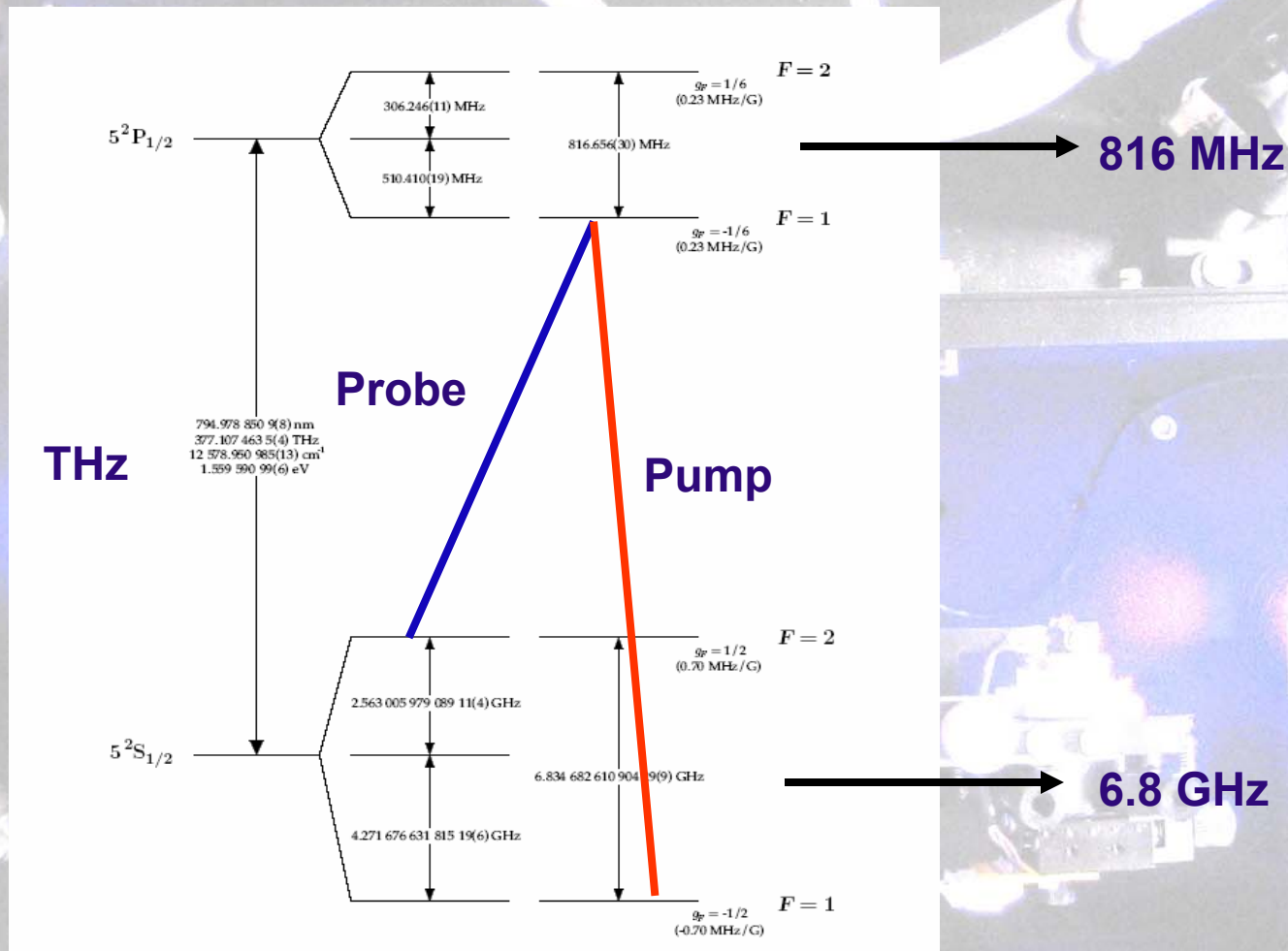
- Reducing the velocity to zero will store the light pulse. (Turning control field off)
- The process can be reversed, thus implementing a quantum memory cell for light.
- This method is good for storing quantum and classical information alike.
- The regenerated pulse should possess exactly the same quantum properties.
- Information is stored in the atoms



D. Phillips, A. Fleischhauer: Phys. Rev. Lett. 86, 783 (2001)

# EIT in an atomic medium

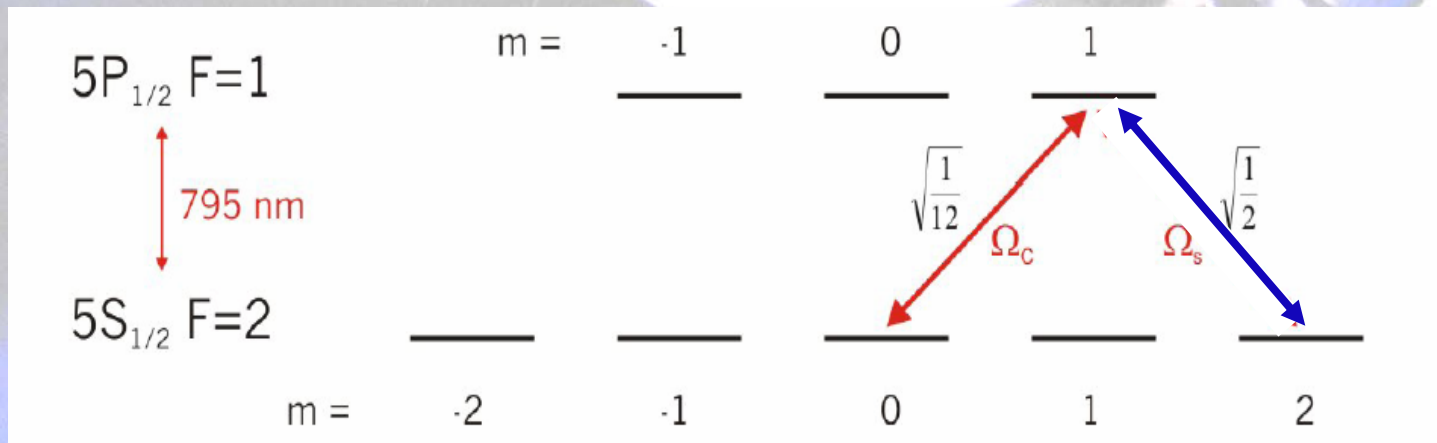
- Atomic Rubidium vapor (D1 line at 795nm)



# EIT in an atomic medium

An alternative scheme:

Zeeman substates of the hyperfine ground states

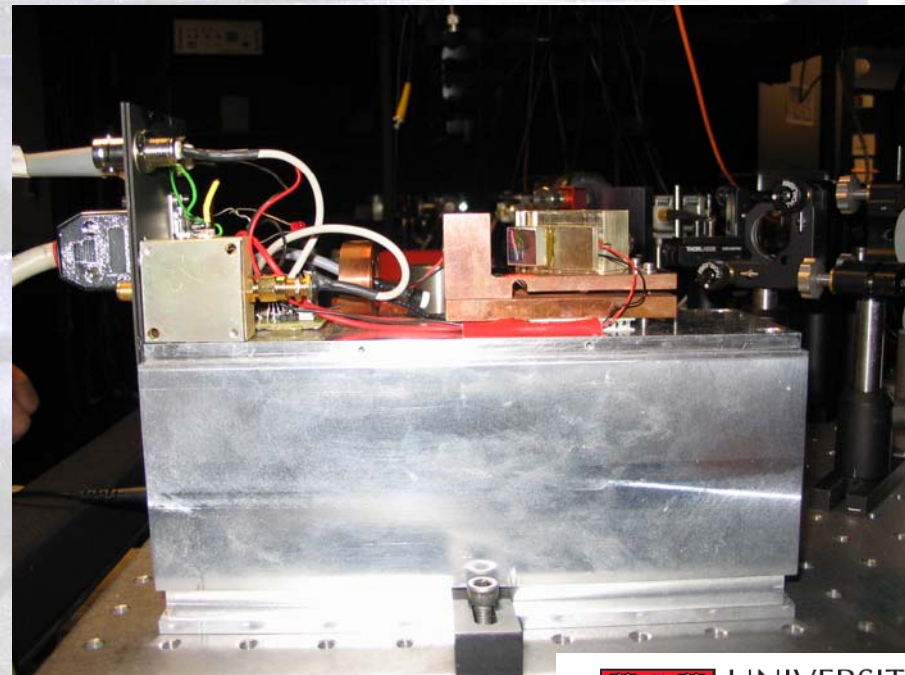
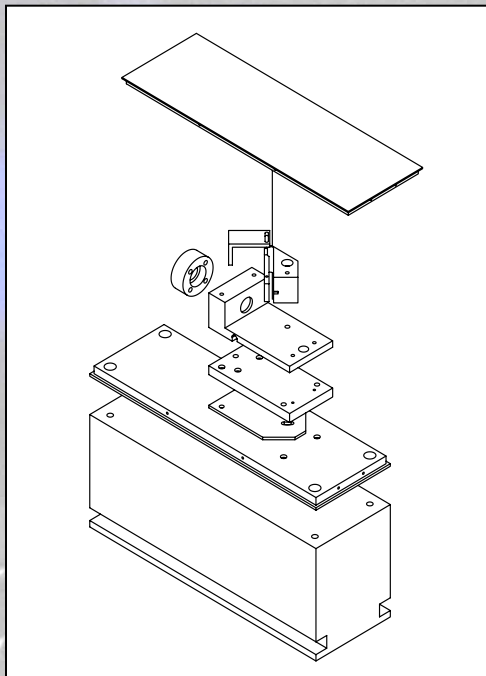
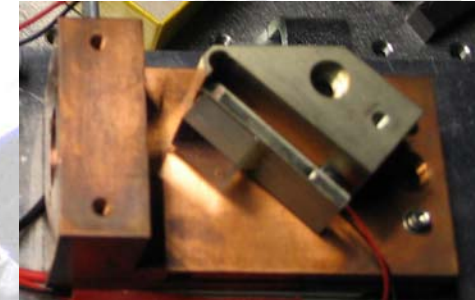


We can use different polarizations of the same laser to observe EIT

Now we go to the real experiments....

# Laser diode system

- Tunable to the rubidium transitions
- Narrow linewidth.
- Keep a constant temperature in the diode
- The electronic control drives the piezo scan over the band of achievable wavelengths.



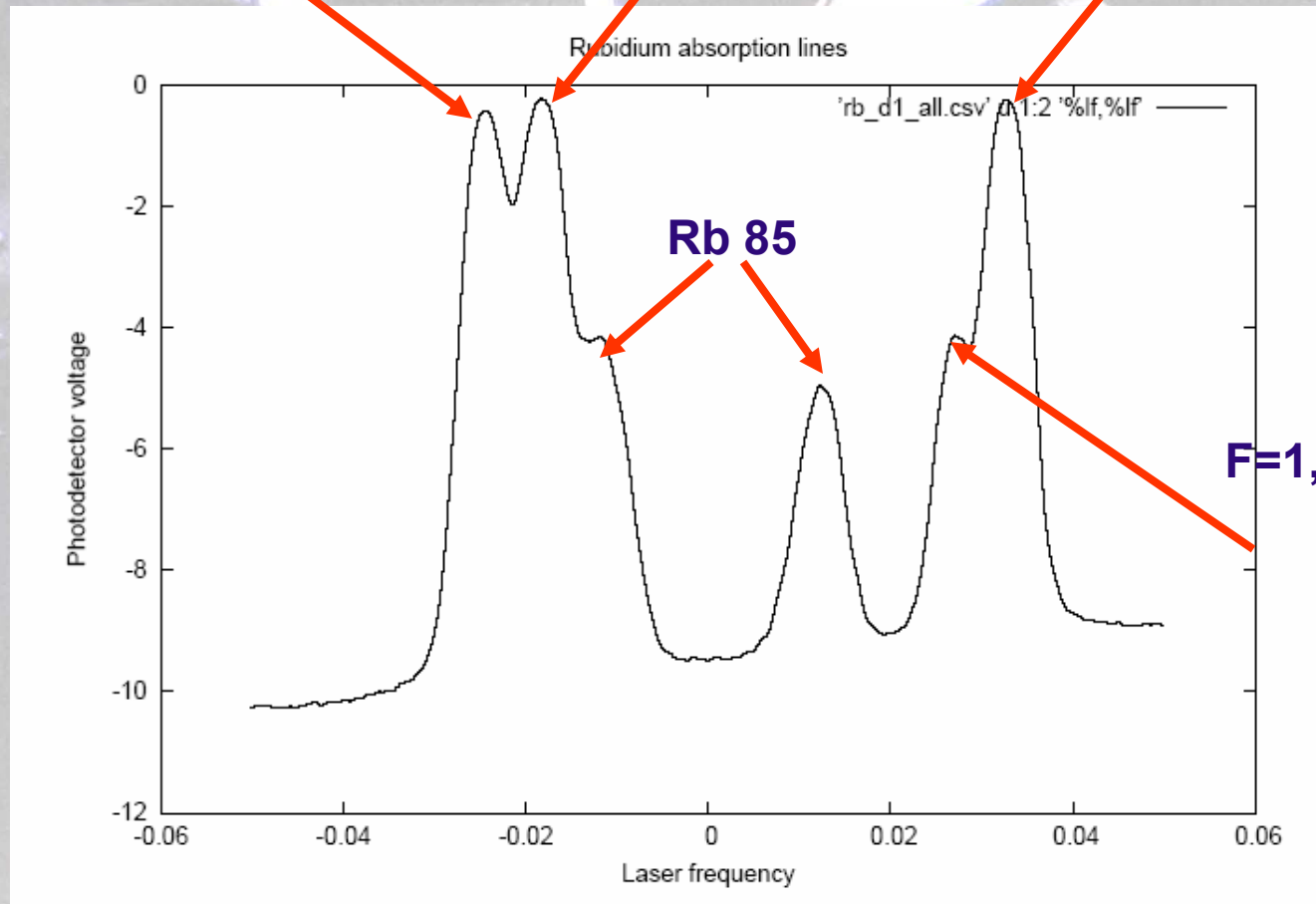


# Rubidium spectroscopy

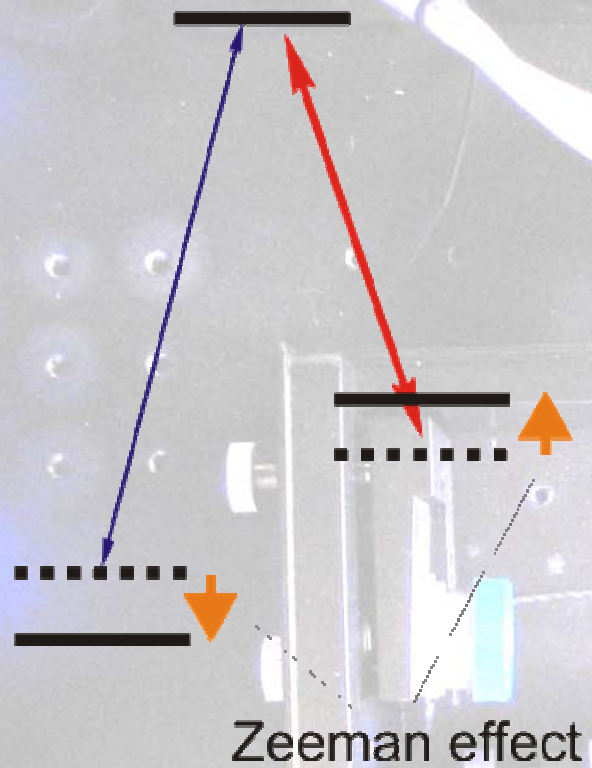
$F=2, F'=1$

$F=2, F'=2$

$F=1, F'=2$



# The effect of magnetic field on EIT



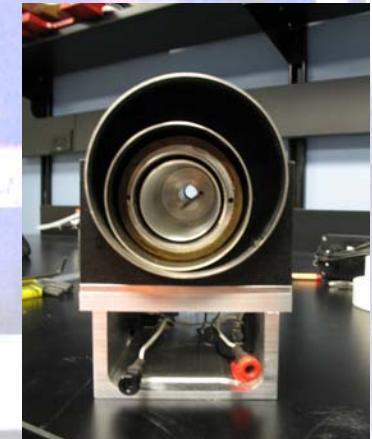
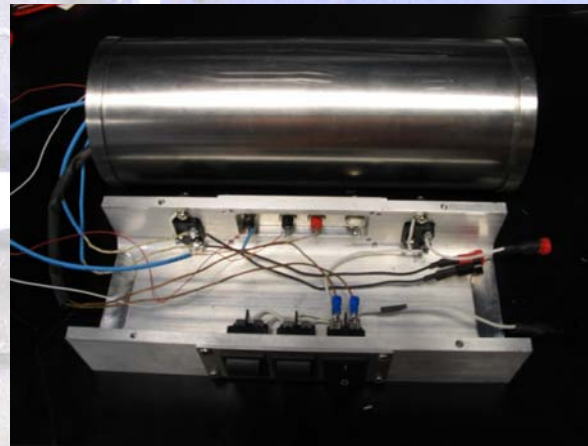
**Magnetic field**

- Zeeman effect
- Levels get shifted
- Two-photon resonance destroyed
- EIT lost

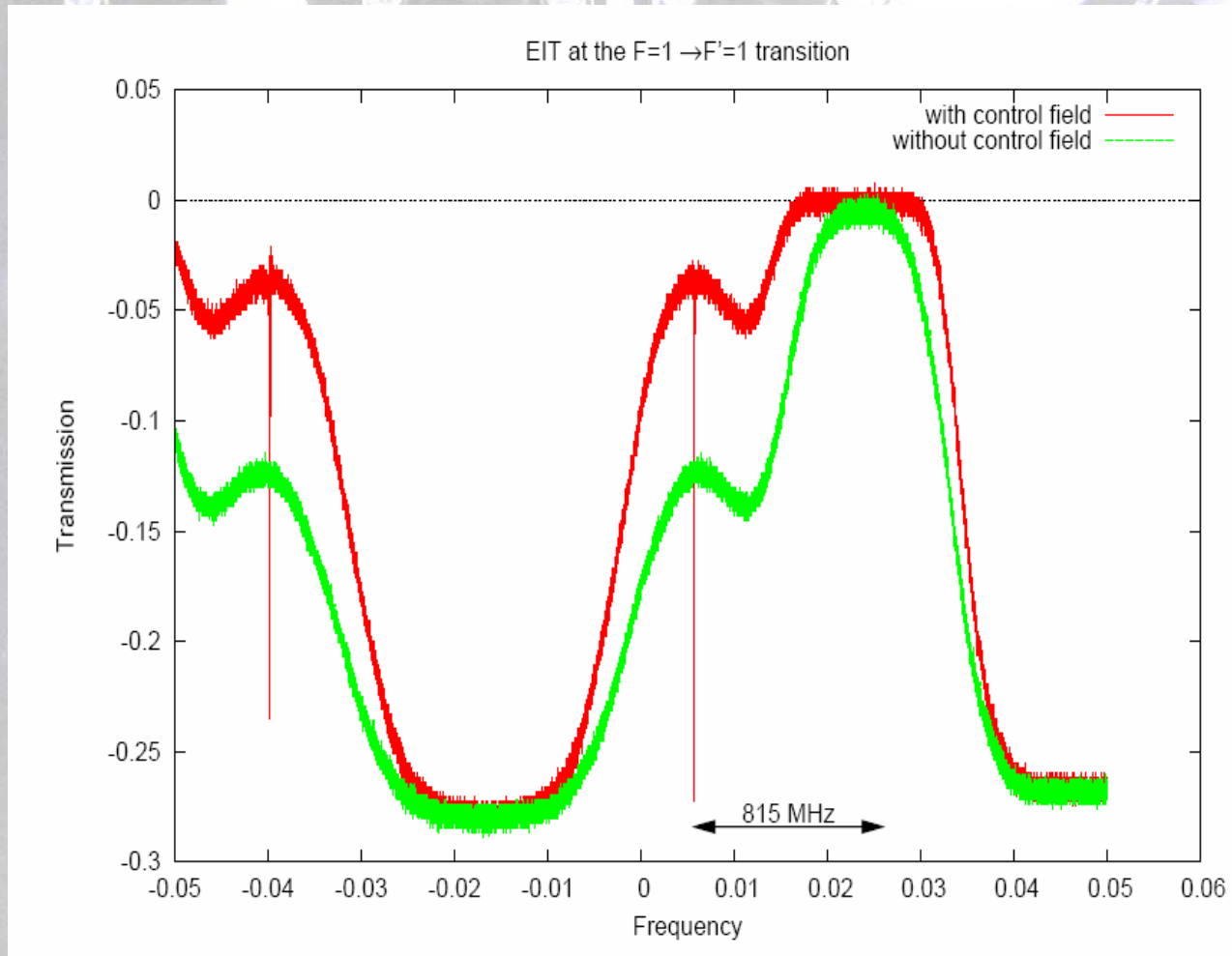
**Magnetic field destroys EIT**

# Interaction zone

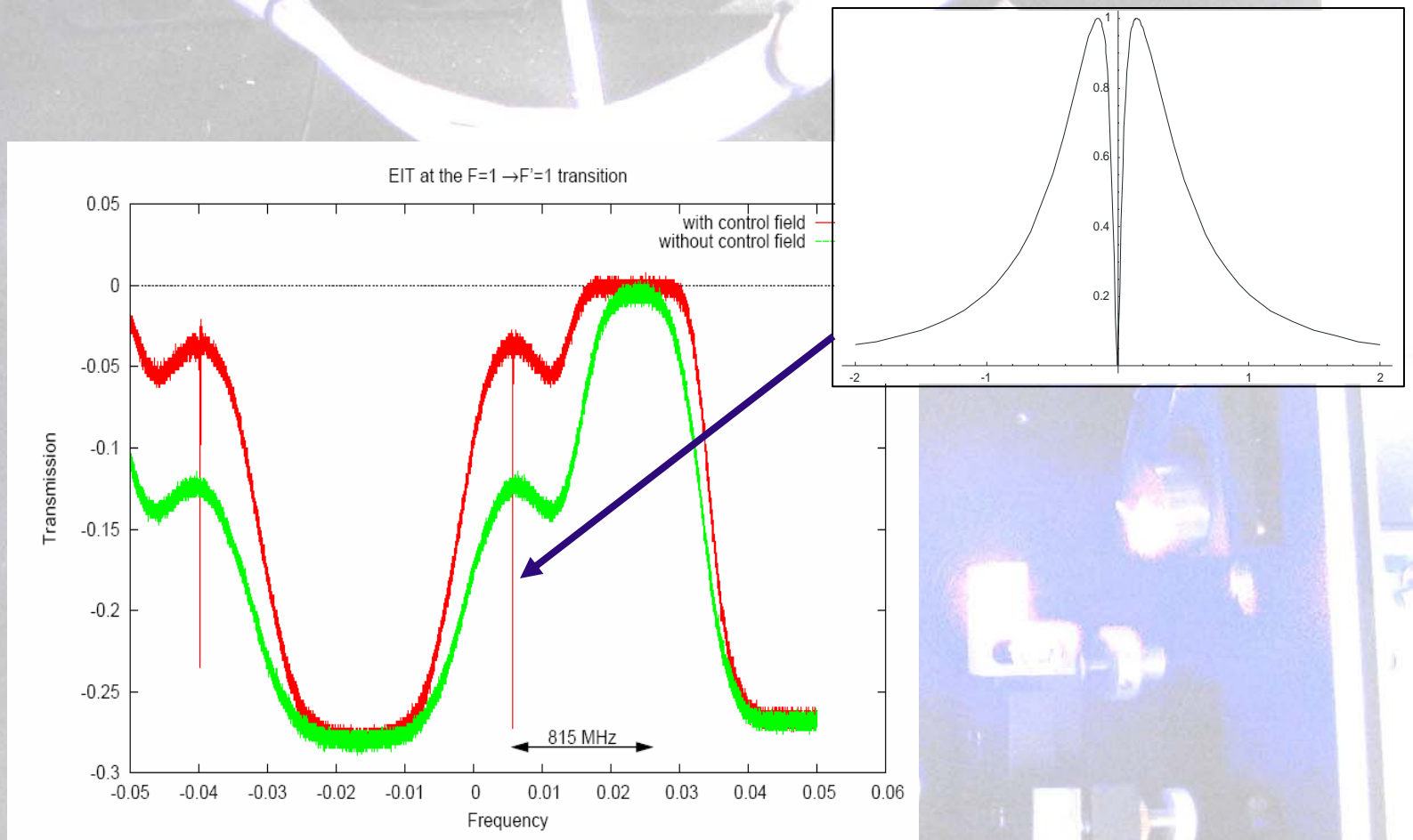
- We have to shield our Rubidium cells.
- Cylinders of Mu-metal are nested together.
- Degaussing procedure is necessary.
- The level of shielding reaches below the micro gauss level.



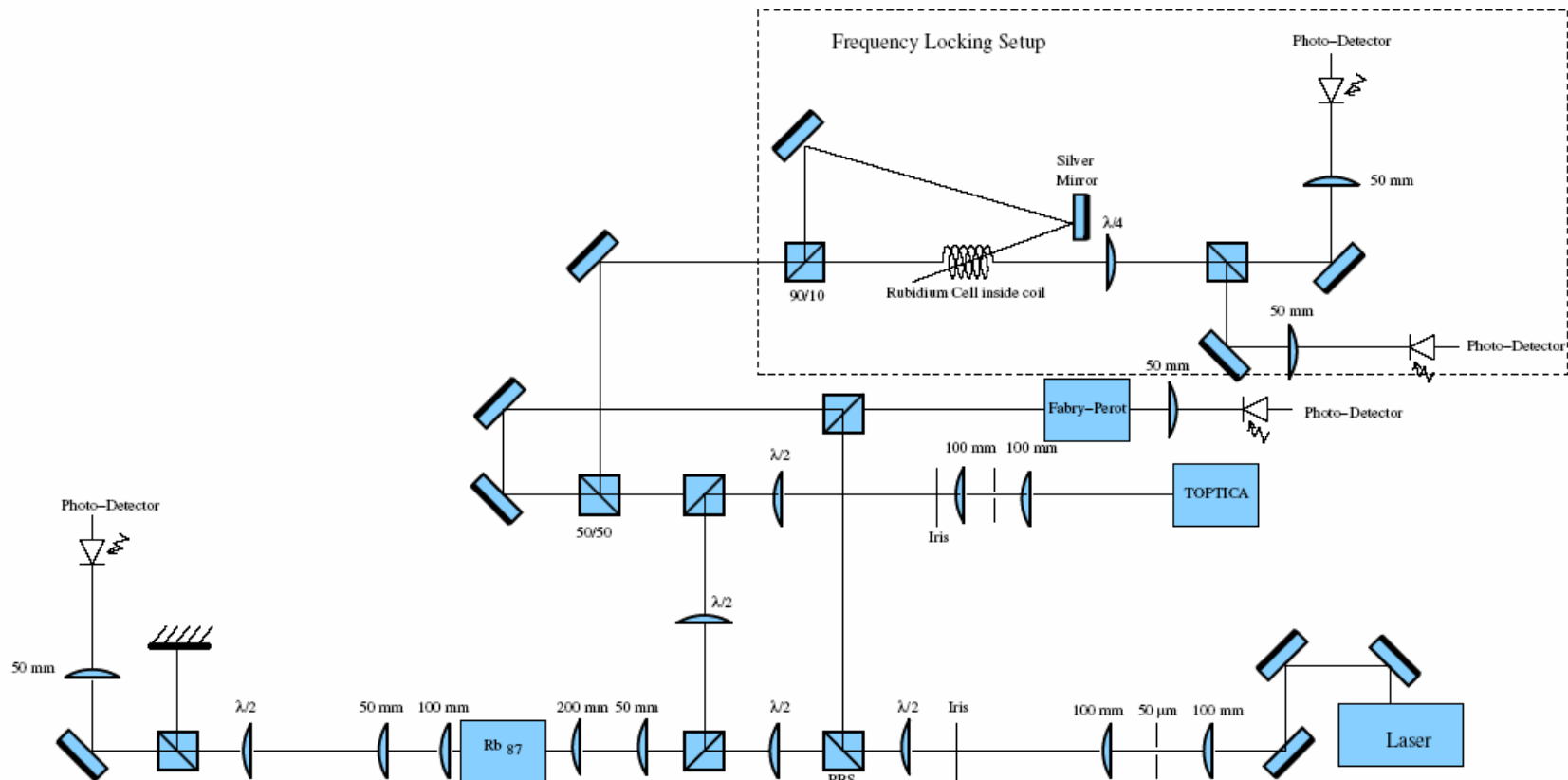
# EIT in the hyperfine levels of Rubidium



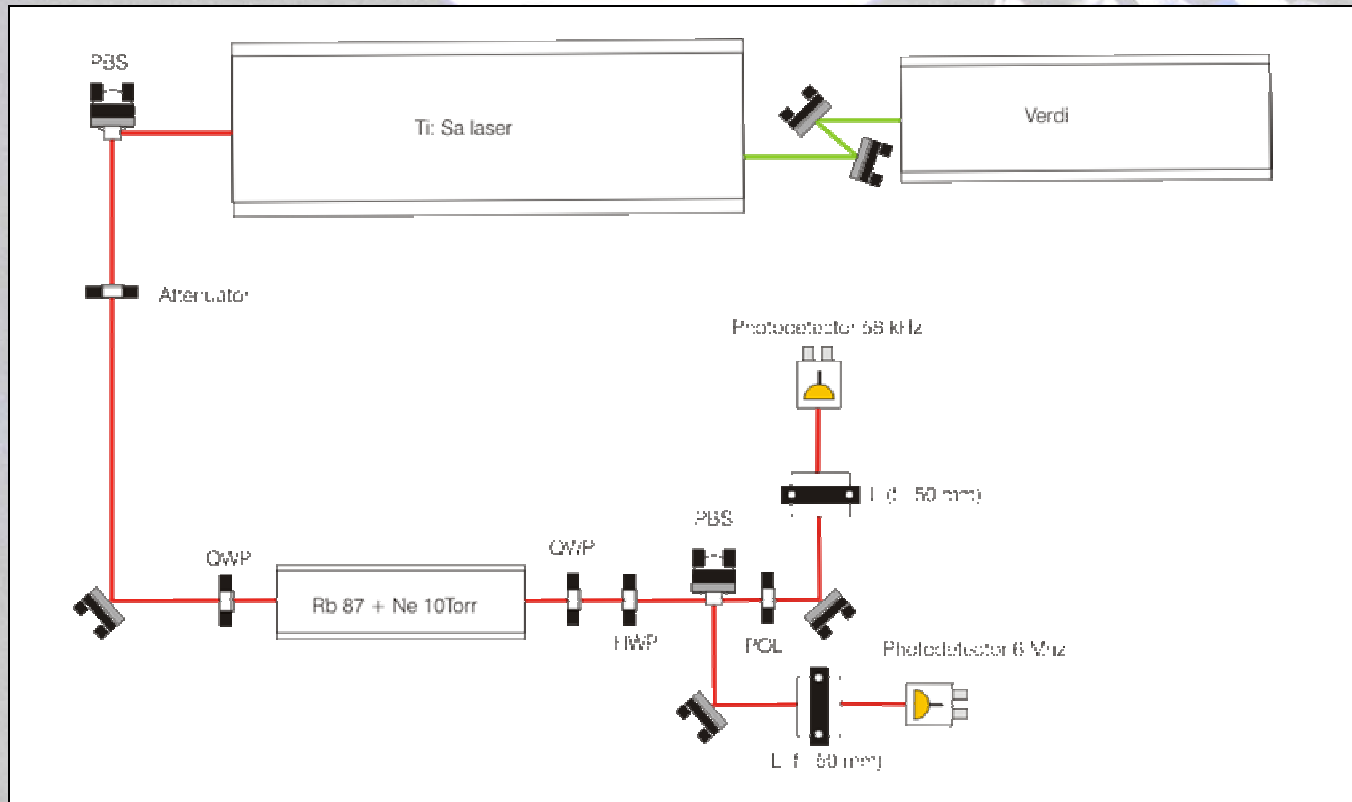
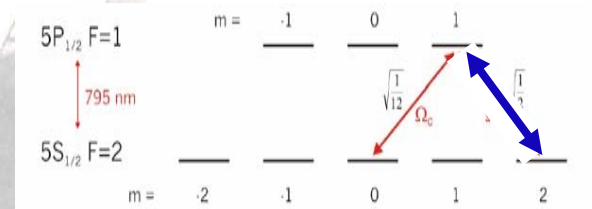
# EIT in the hyperfine levels of Rubidium



# EIT in the hyperfine levels of Rubidium



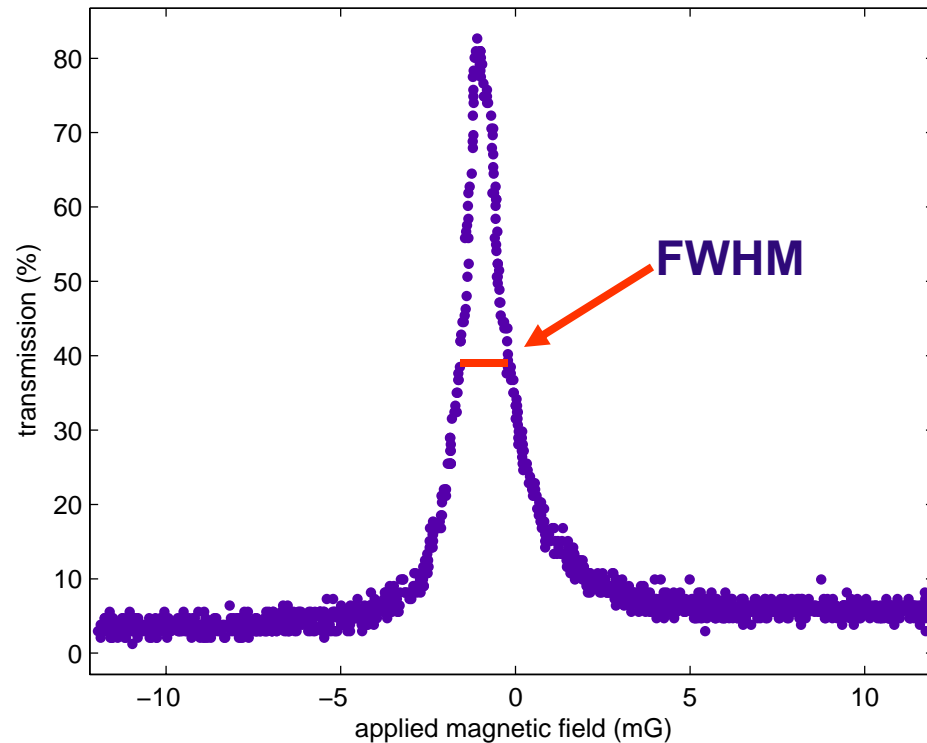
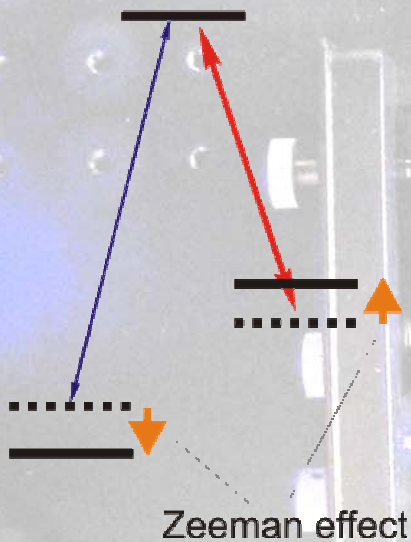
# EIT in Zeeman sublevels



# Measuring the EIT width (Method I)

We want to measure the width of the EIT transparency window because it contains information on the maximum storage time.

We detuned the two-photon resonance using magnetic field.

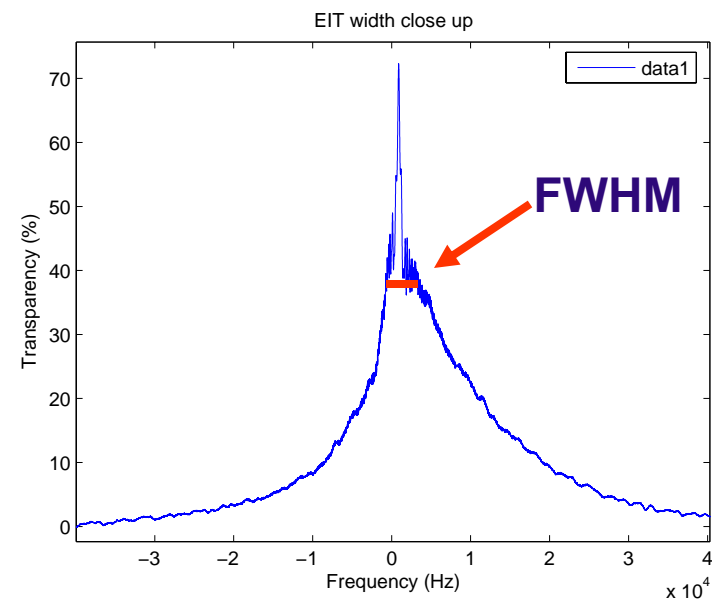
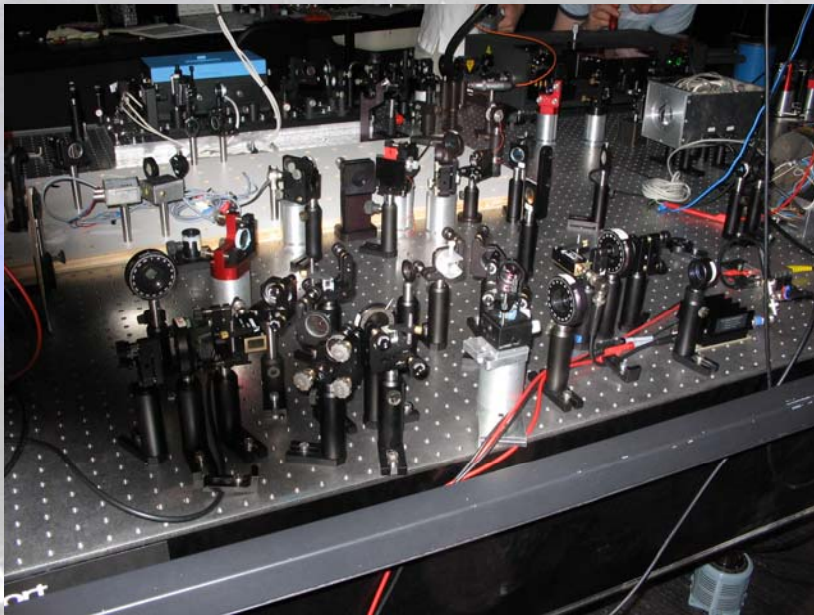


2 mG FWHM  $\rightarrow$  2.5 kHz EIT width



# Measuring the EIT width (Method II)

We tuned control laser vs. signal laser using an Acousto-optical modulator.  
Our EIT width is about 10kHz FWHM, that means our laser should be stable at that level



# Measuring the EIT width (Method III)

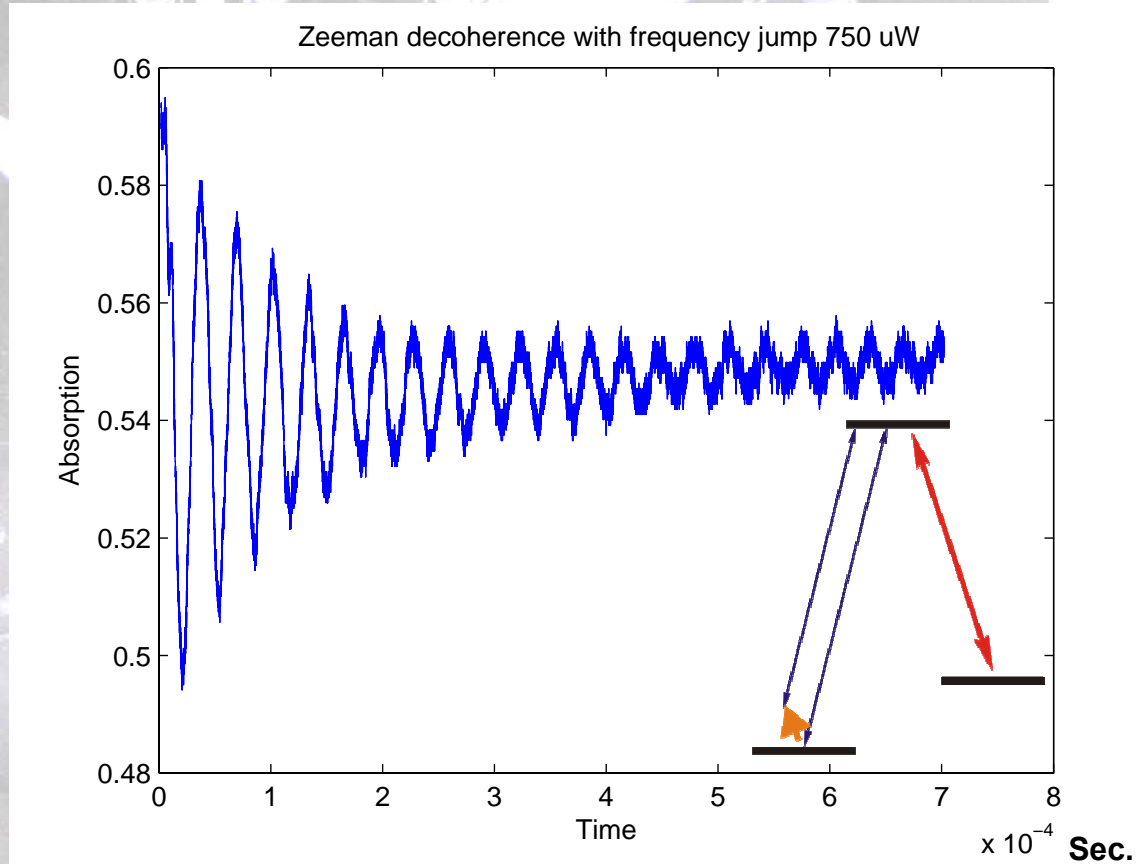
Observe EIT → Ground state coherence is created

Shift laser from two-photon resonance

Ground state coherence will oscillate with respect to the lasers.


Oscillations will damp when ground states decohere

Decoherence time can be measured



S.Park, H. Cho: Phys. Rev. A 69 023806 (2004)

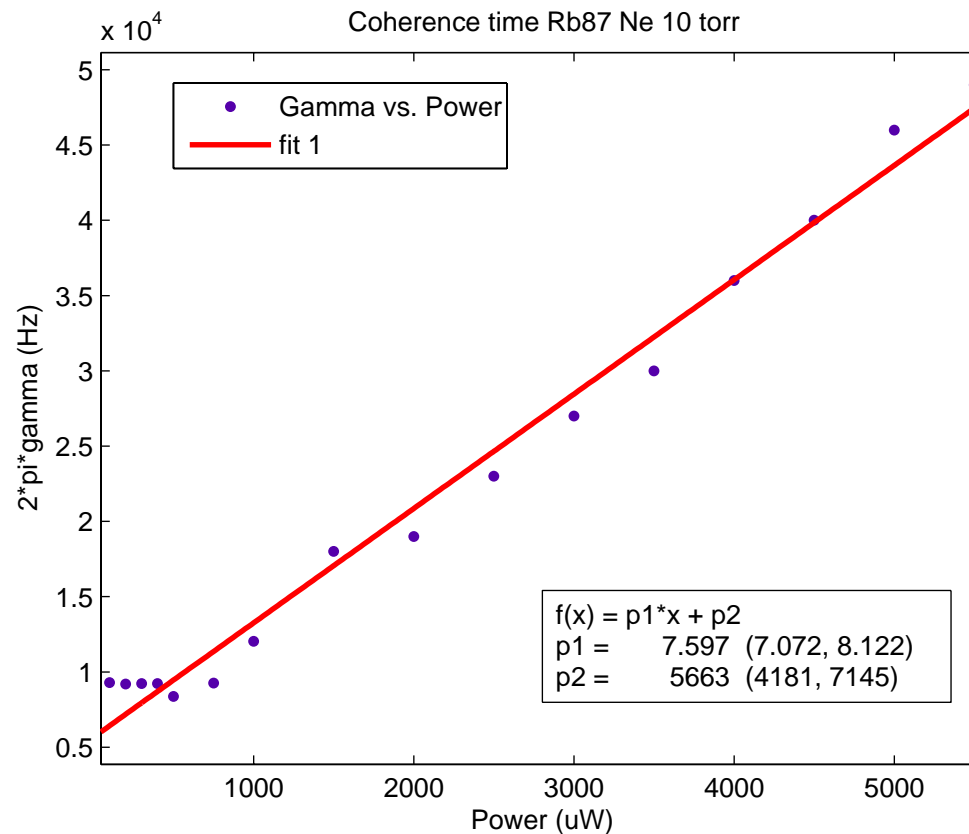
# Ground state coherence lifetime: power dependence

Ground state  
decoherence time   
Inverse of the EIT  
width

Large fields → power  
broadening

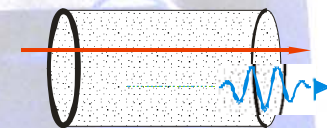
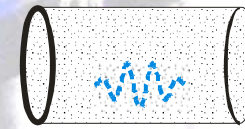
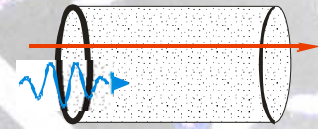
Small fields → ground  
state decoherence  
measurement

- The decoherence  
time is the ultimate  
limit to store light.



# Storing light by EIT

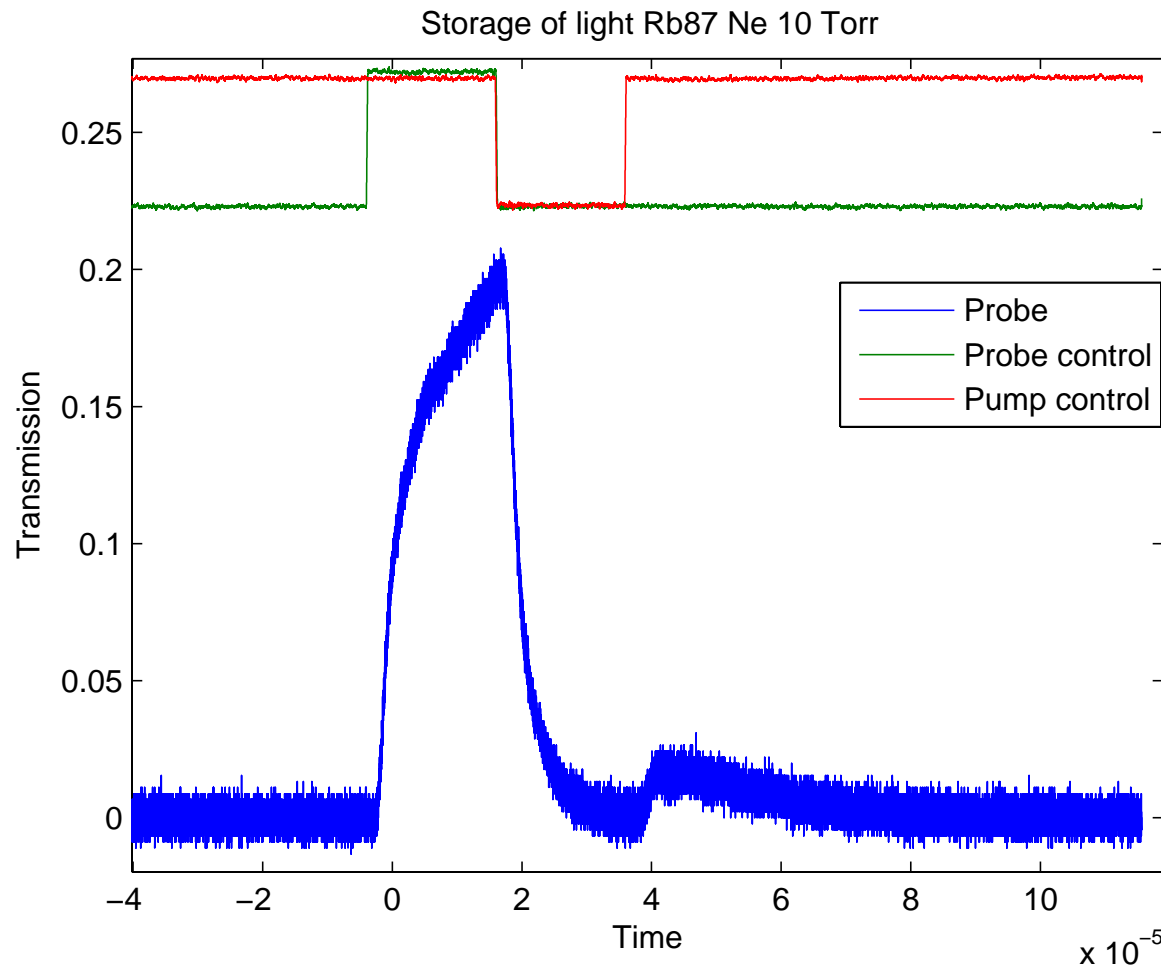
- 1 Couple a pulse of light into an EIT medium
- 2 Turn control field off adiabatically  
→ quantum state of the signal pulse stored as a collective atomic excitation
- 3 Turn control field back on  
→ optical pulse released in the original quantum state



## Storing light by EIT (Technical challenges)

- We need pulses in the probe field. Switch via Acousto-optical modulators
- The group velocity is proportional to the power of the pump field.
- We need to smoothly switch off the pump field to reach the zero of the group velocity.

# Storage of light (Preliminary results)



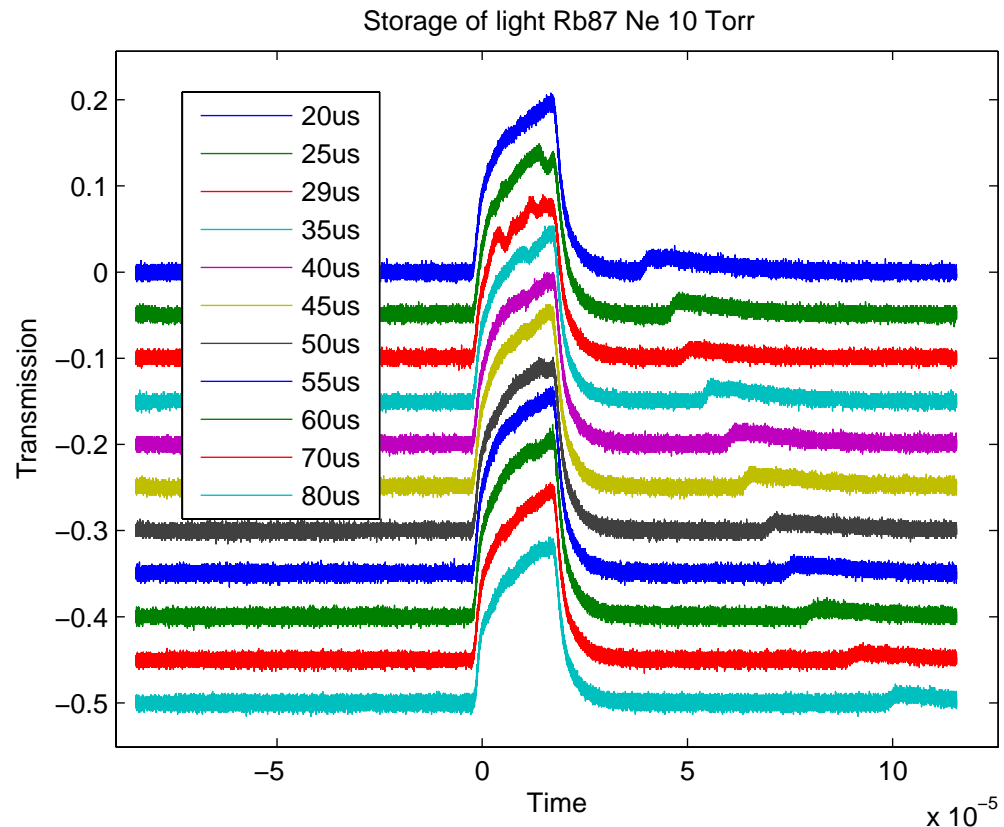
**Voila!!**  
**Stored**  
**light**

# Storage of light (Preliminary results)

Storage of light up to 80  $\mu\text{s}$

Storage can be improved by using Gaussian pulses

80  $\mu\text{s}$  = 240 km in free space



# Ultimate goal: Storage of non-classical light

## Challenge:

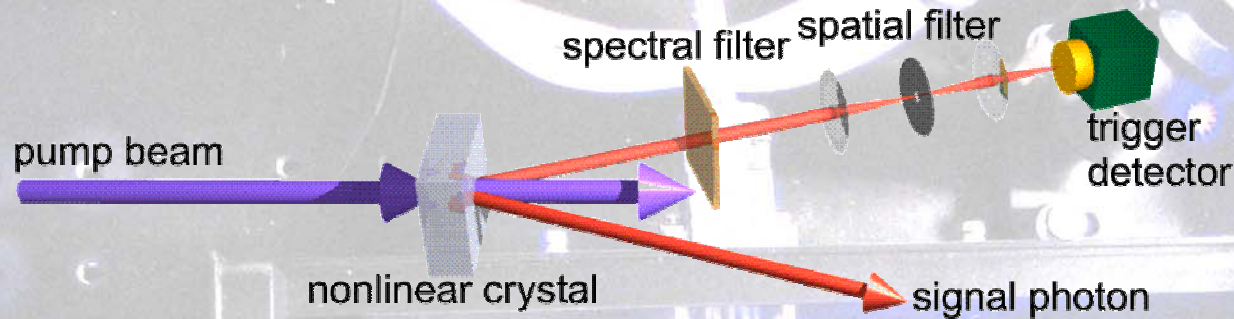
- We require a narrowband and tunable source of non classical light for experiments with atoms

## Method:

- Parametric down conversion in a resonant cavity



# Parametric down-conversion

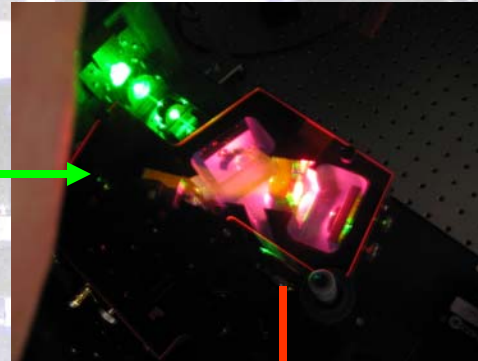


- “Red” photons are always born in pairs
- Photon detection in one emission channel  
→ there must be a photon in the other channel as well (Heralded photons)

# Single photon source

Pump laser 10W (560nm)

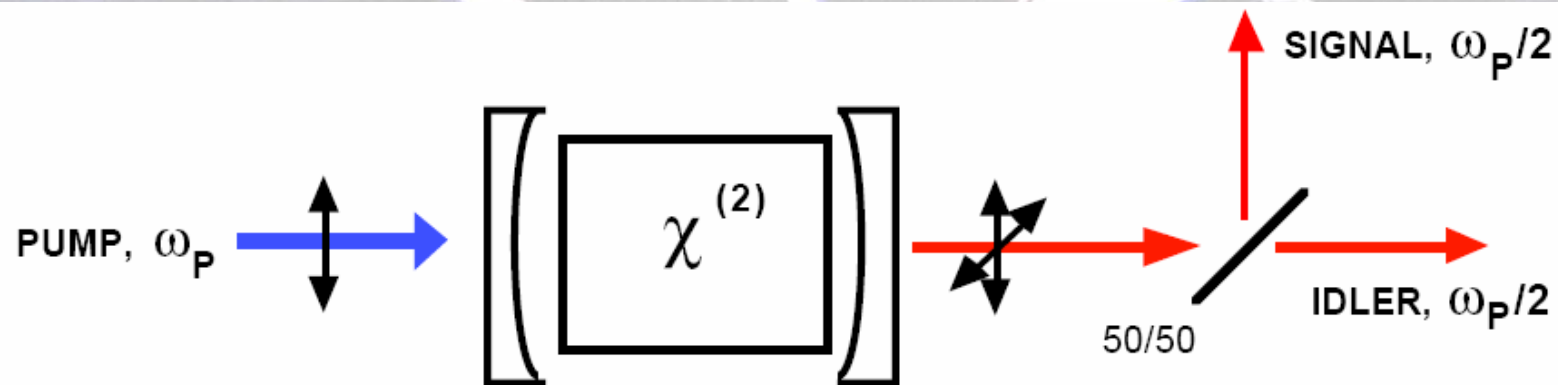
Ti:Sapphire laser 1.8 W (795nm)



OPO

Frequency doubler 700 mW (397.5nm)

# Optical parametric oscillator



- This cavity must be stabilized so that the signal photon wavelength is exactly the same as that of the desired atomic transition.
- One photon in a pair will be detected while the other one will be stored in an EIT medium.
- We will verify whether the EIT storage preserves the single-photon state.

# Summary

- Building and characterization of a tunable laser system
- Production of EIT configurations in Rb 87 atoms
- Design of a magnetic shielding together with temperature stabilization for our cells.
- Measurements of EIT width by different methods
- Measurements of decoherence times of ground states
- Preliminary results on storage of light
- Overview of the single photon source

## Future work

- Implementation of an OPO pumped with light coming from a frequency-doubling cavity to be used as a source of non-classical states of light.
- Design and characterization of a filtering system for narrow band single-photons.
- Storage and retrieval of single photons in Rubidium vapor.

**Thanks to:**

**Lab visit at 5 PM!!!**

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German Academic Exchange Service



**National Research  
Council Canada**

**Conseil national  
de recherches Canada**



**Eden Figueroa CS-QIC 2005**

