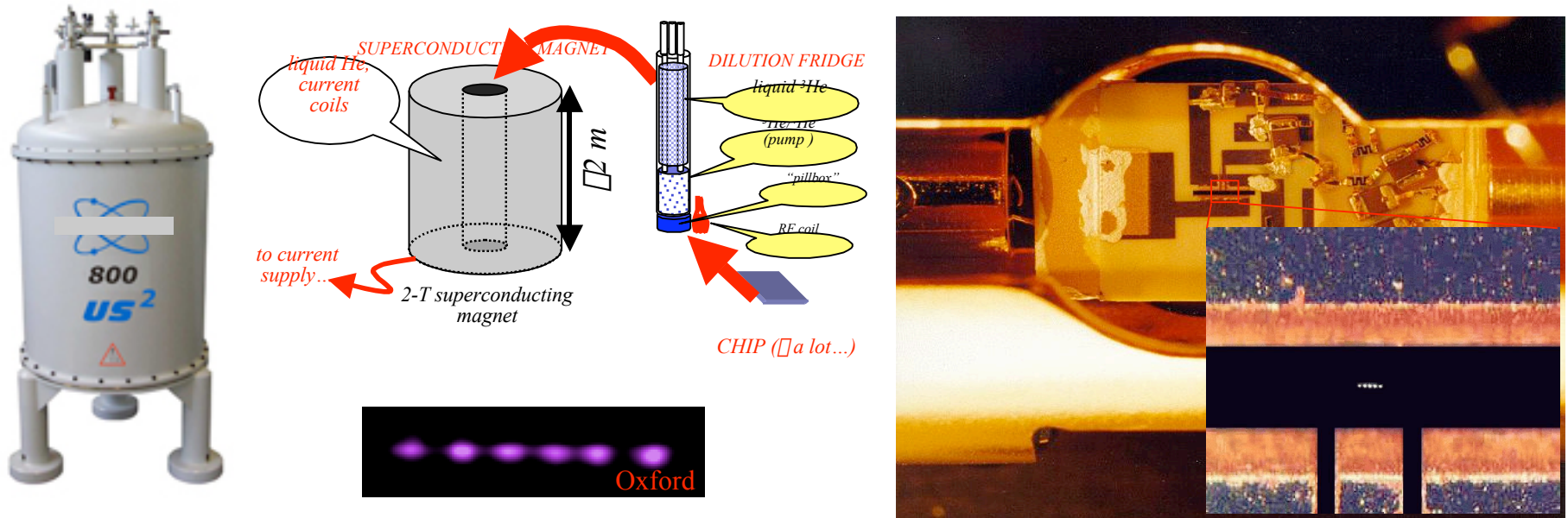


# Physical "Implementations" of Quantum Computing



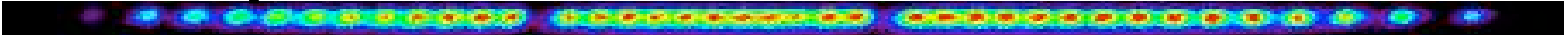
**PIMS/MITACS** summer school– June, 2003

**Brian King**

Dept. Physics and Astronomy, McMaster University

[http://physserv.mcmaster.ca/~kingb/King\\_B\\_h.html](http://physserv.mcmaster.ca/~kingb/King_B_h.html)

# Frontspiece:



*“Quantum mechanics  
does not occur in a  
Hilbert space – it occurs  
in a laboratory!”*

*- Asher Peres*

# Outline:

- propaganda ✓

## Part I:

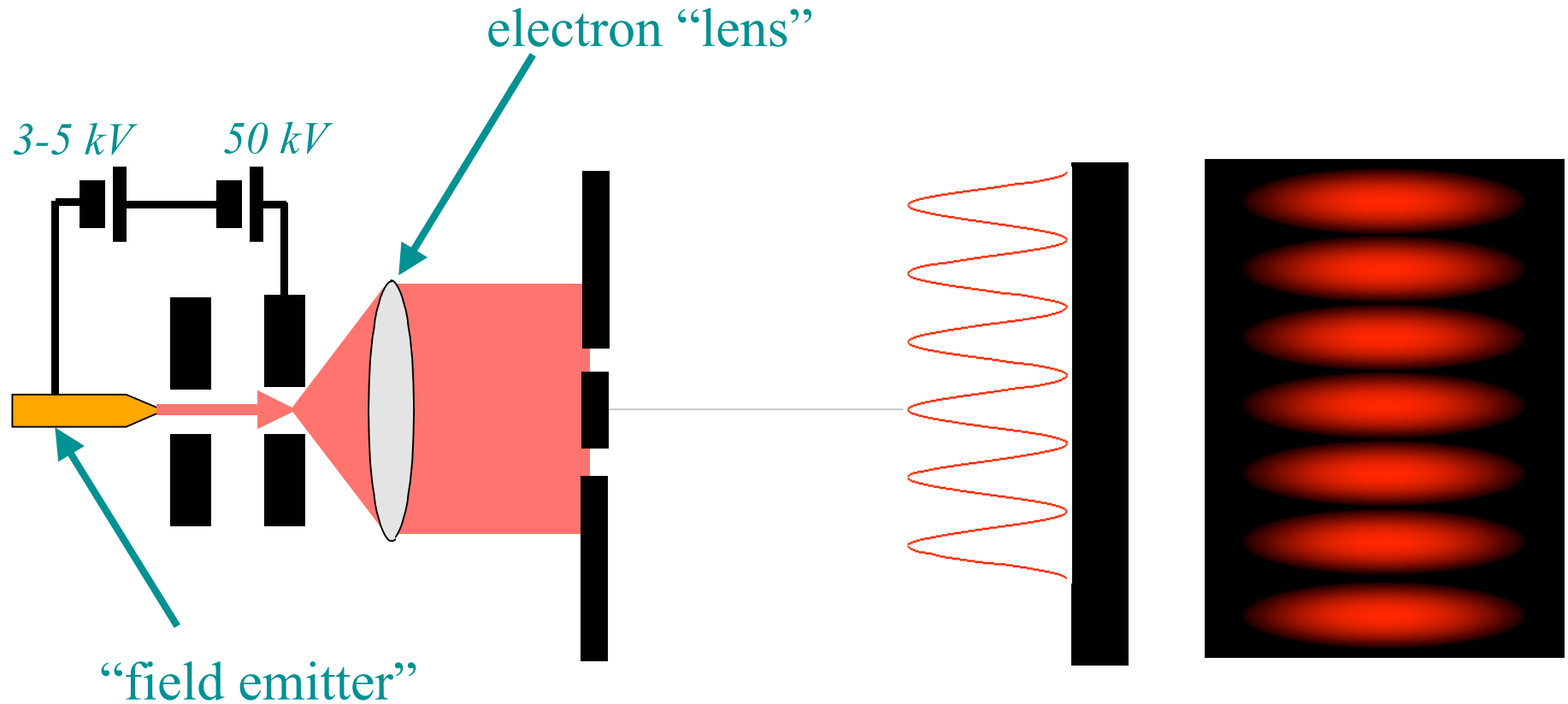
- quantum behaviour
- quantum dynamics
- “real” qubits (by way of atoms...)
- physical requirements
- “rogue’s gallery” of architecture proposals

## Part II:

- ion traps
- atomic qubits, again
- logic gates
- the future (?...)

# Electrons and holes!:

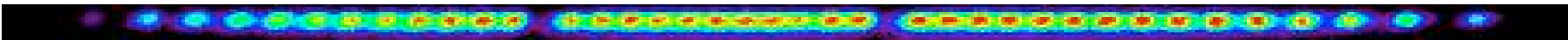
A Tonomura, *et al.*, 1989 Am. J. Phys. **57**, 117-120 (89); L Marton, *et al.*, Phys. Rev. **90**, 490-491 (53); C Jönsson Zeit. Phys. **161**, 454-474 (61)



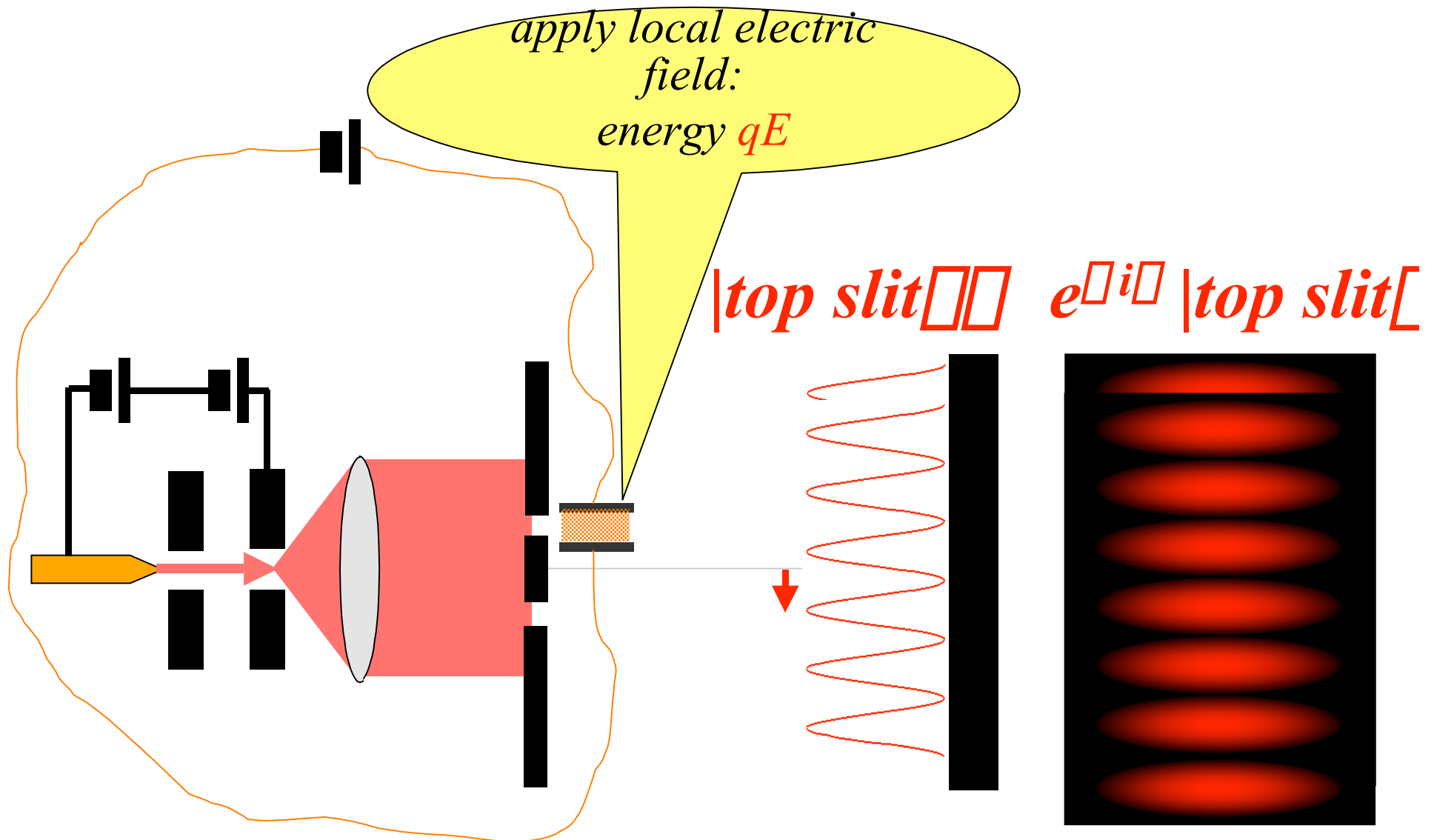
*state of electron*  $\sim |top\ slit\rangle + |bottom\ slit\rangle$



# Electrons and holes!:



- locally change part of the electron's energy...



# Quantum dynamics:

- the Schrödinger equation:
  - tells us how the quantum state *changes in time*

$$i \hbar \frac{d}{dt} |\psi\rangle = \hat{H} |\psi\rangle$$

*quantum state  
(state vector)*

***Hamiltonian  
operator***

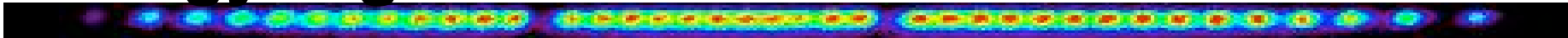
• *energy of  
quantum state*

- energy *eigenstates*:

$$\hat{H} |\psi_{E,n}\rangle = E |\psi_{E,n}\rangle \quad E \text{ is the energy of } |\psi\rangle$$

- *state vector is unchanged (up to a constant...)*

# Energy Eigenstates:



- *make solving the Schrödinger eq. easy!*

$$H |\psi_{E,n}\rangle = E |\psi_{E,n}\rangle$$

$$i \hbar \frac{d}{dt} |\psi_{E,n}\rangle = E |\psi_{E,n}\rangle$$

$$|\psi_{E,n}\rangle(t) = e^{-iE_n t / \hbar} |\psi_{E,n}\rangle(t=0)$$

- *linear vector space*  $\psi$  solve other states by superposing solutions to energy eigenstates

# Energy Eigenstates:

aside: how do we figure out energy eigenstates?

- usually amounts to solving spatial PDE:

spatial distribution of  
quantum state  $|\psi\rangle$

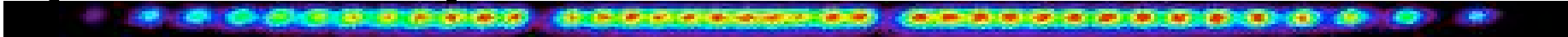
$$\psi(x) = \langle x | \psi \rangle$$

$$\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x)}{\partial x^2} = (E - U) \psi(x)$$

total energy of  
quantum state  $|\psi\rangle$

potential energy (gives  
the forces)

# qubits vs. “qubits”:



- *so why do we care?!? ...*

## qubits:

- two-level quantum system
- basis states  $|0\rangle$  and  $|1\rangle$
- **superpositions**  
 $\alpha |0\rangle + \beta |1\rangle$

## “qubits” (real world):

- in a non-empty universe, multi-level quantum systems
- **superpositions**  
 $\alpha |0\rangle + e^{iE_{01}t/\hbar} \beta |1\rangle$ 
  - must track phases!  
 $\hbar$  stable oscillator
- **decoherence!**

# Building Quantum Computers:

## To build a “real” quantum computer, need:

### 1. qubits

- two-level quantum systems (or effective ones...)
- *superpositions* □ isolated from outside world
- confined, characterizable, scalable

### 2. preparation

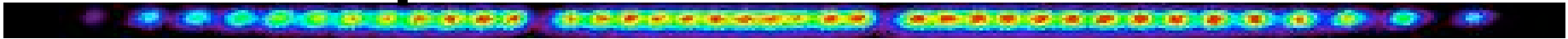
- prepare computer in standard start state

### 3. read-out

### 4. logic gates

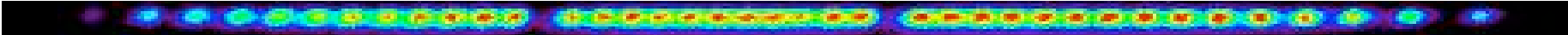
- controllable interactions with outside world!
- single- and two-qubits gate sufficient (not nec.!).

# Atoms as qubits

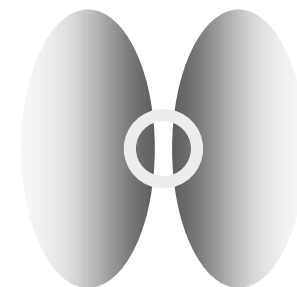
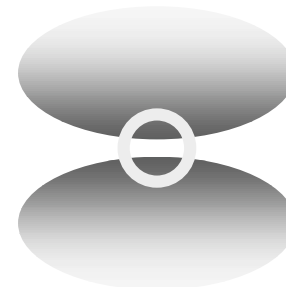
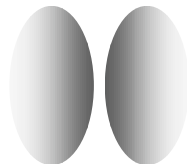
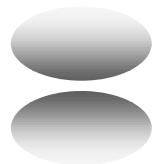
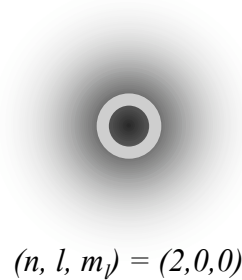
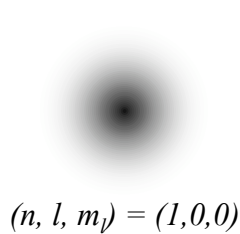


- why atoms?
  1. (my) familiarity...
  2. atoms are *the* standard for quantum superpositions
    - definition of the second:
      - phase evolution of superposition of atomic levels in Cesium
      - accuracy and stability  $\sim 1/10^{16}$

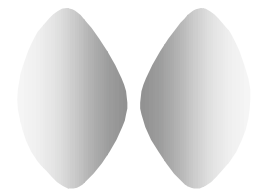
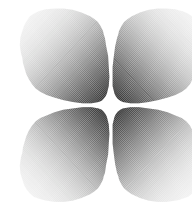
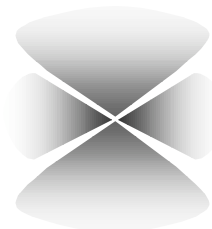
# Quantum atoms:



*atoms: + charged nucleus, □ charged electrons*

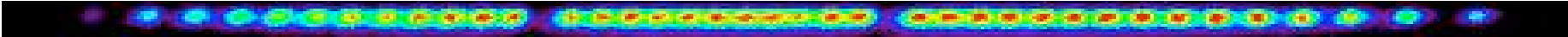


*n ~ h...  
l ~ ...  
m<sub>l</sub> ~ prob. it's rotating about the direction you're looking...*

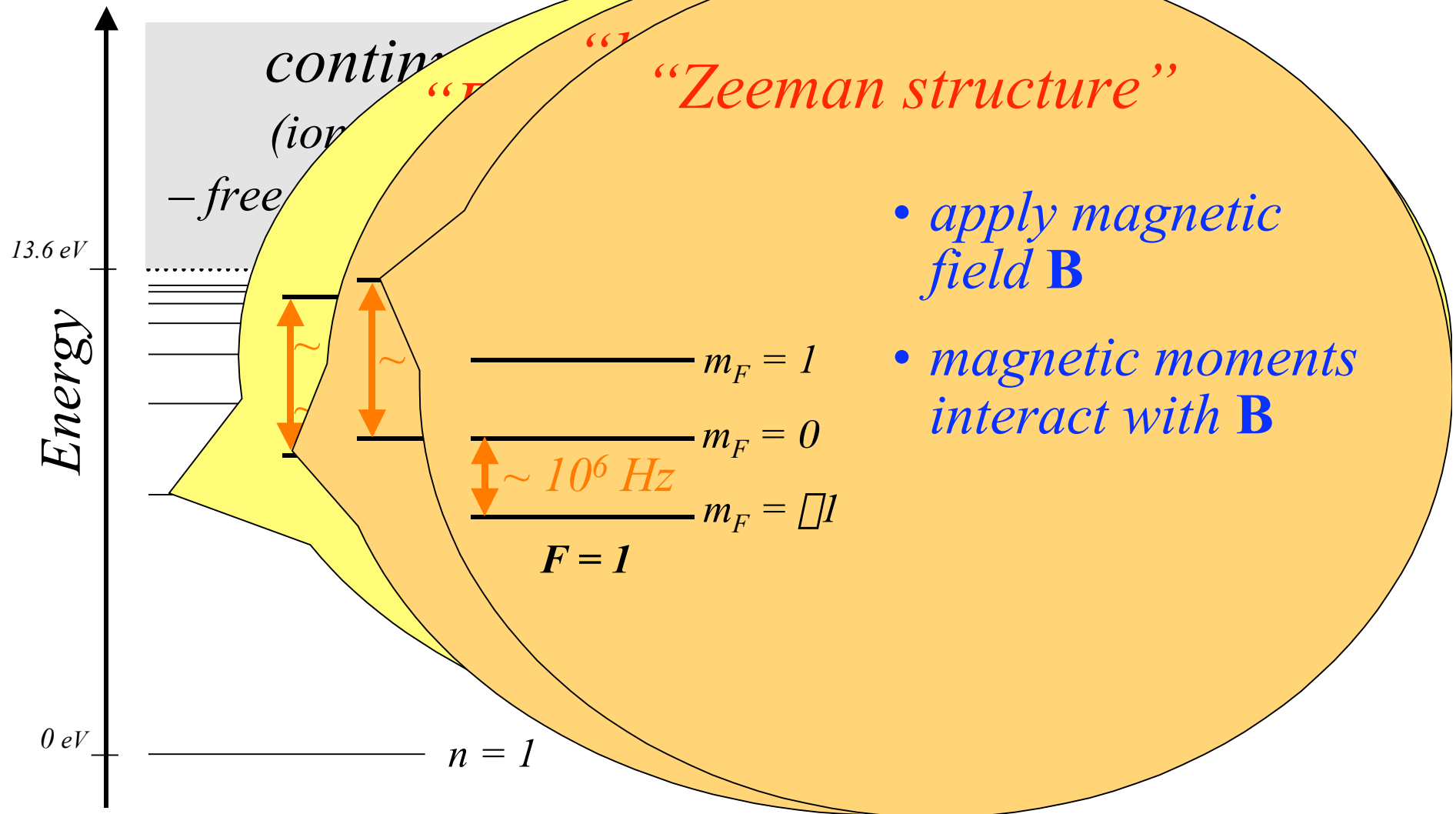




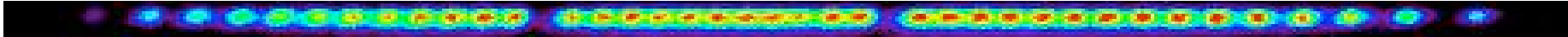
# Atoms:



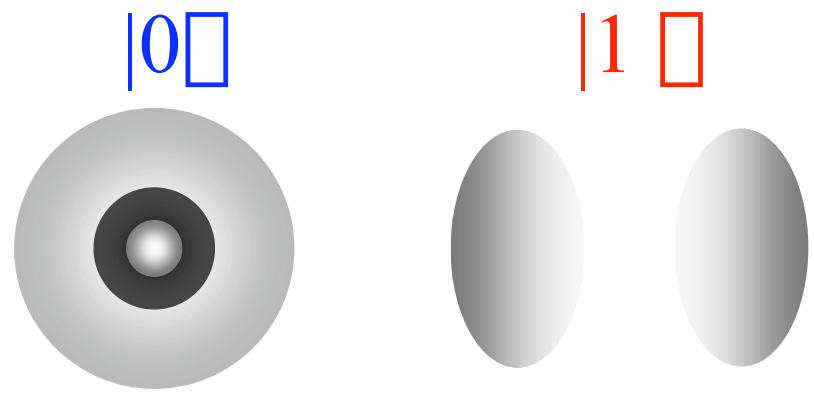
- energy-level diagram
  - e.g. atom with a single electron (e.g. hydrogen)



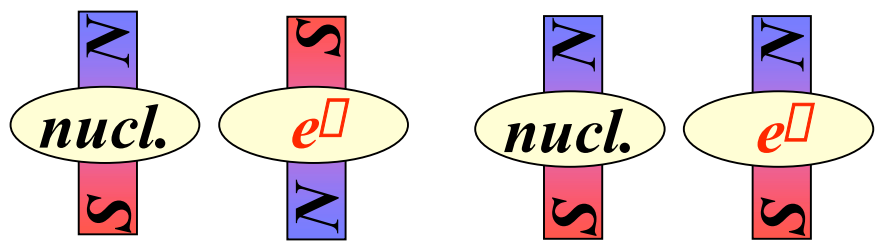
# Atoms:



- **qubits:** isolate two energy levels
  - e.g. lowest-energy level (thermodynamics)
  - + another level you can get to with some coupling



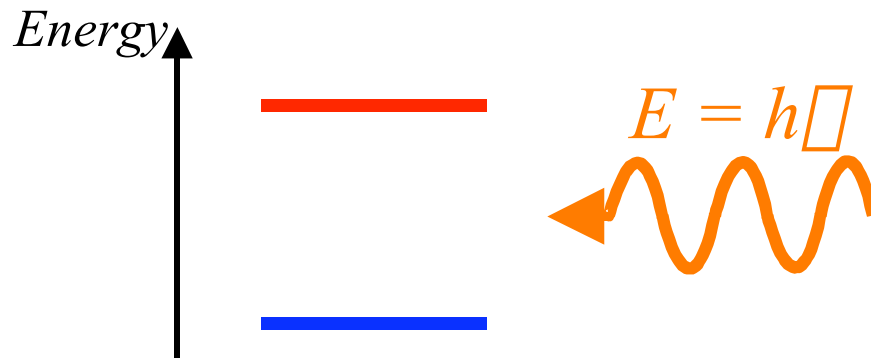
1. long-lived electronic levels ( $T_1 > 1 \text{ ms}$ )



2. long-lived ground-state hyperfine levels ( $T_1 > 10,000 \text{ y.}$ )

# Single-qubit logic gates:

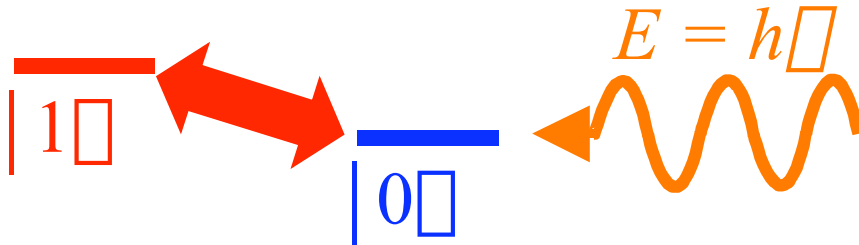
- analog of classical “*NOT*” gate - expanded!
- equivalent to preparation of arbitrary qubit state
- **apply laser/microwaves!**



- absorption  $\uparrow$
- stimulated emission  $\square$
- ~~spontaneous emission~~
  - (long-lived ex. state...)

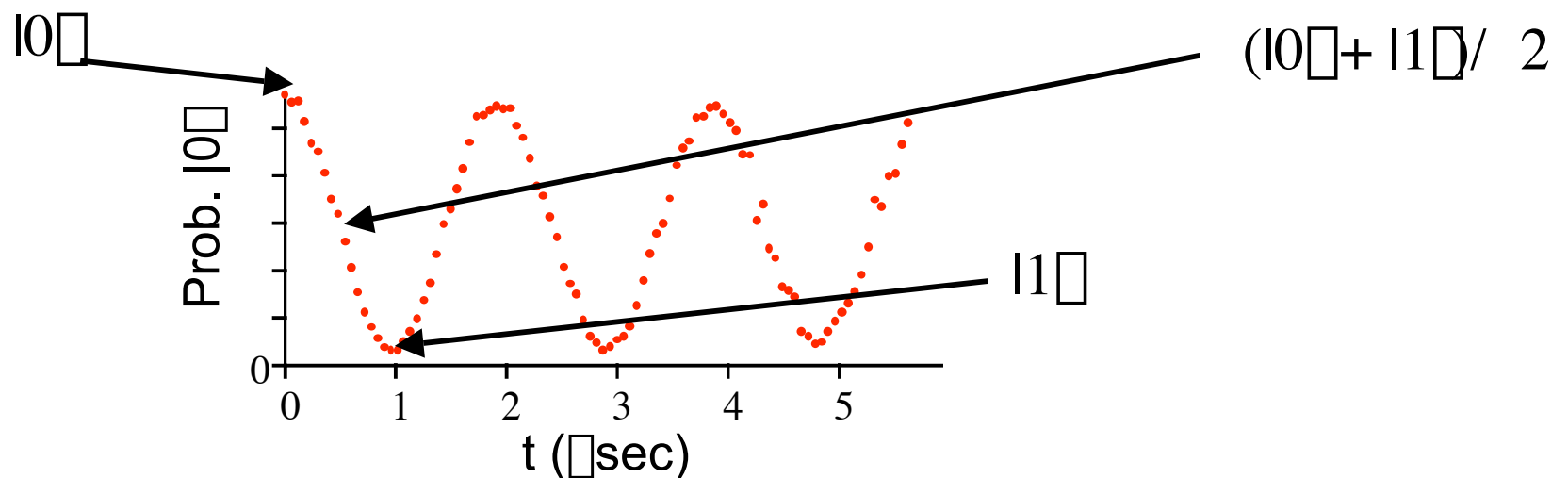
- **frequency must (approximately\*) match the energy “gap” between the energy levels**

# Single-qubit logic gate:



- classical: random hopping (either/or)
- quantum: flow of wave-function with time

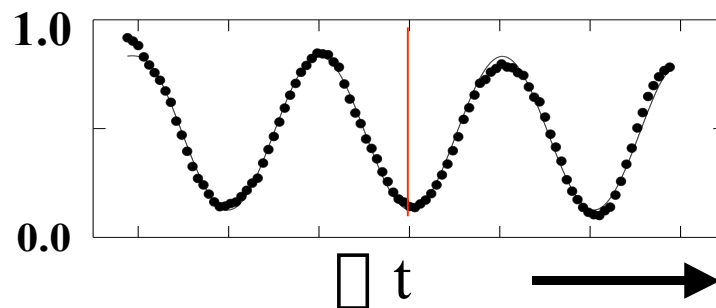
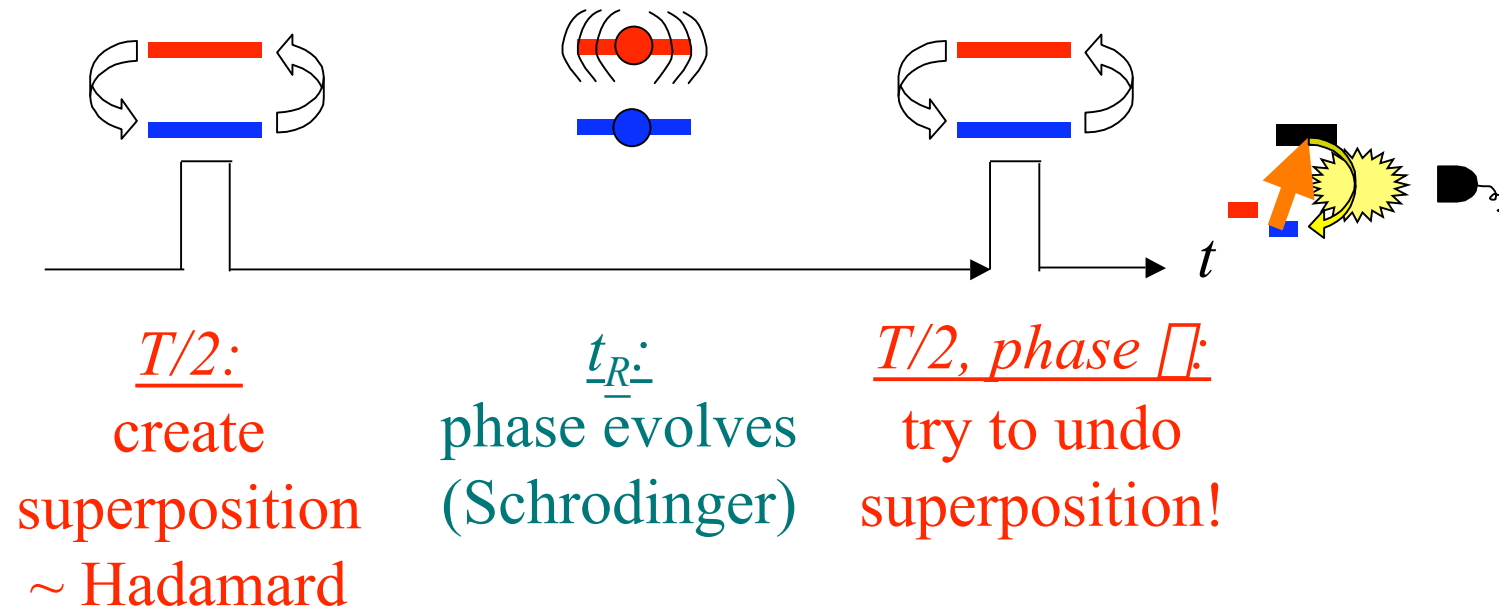
$$E_0 \quad E_1$$



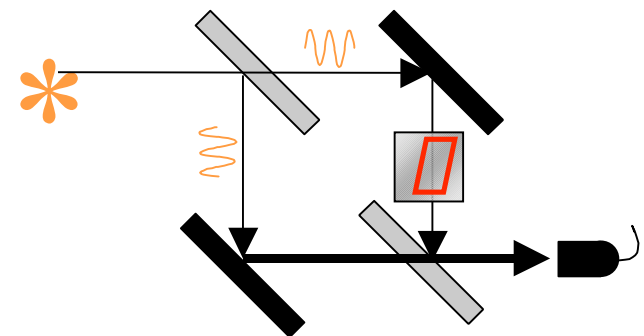
# Nobel Sidebar - Ramsey's expt.:

N. F. Ramsey, Rev. Mod. Phys. **62**, 541-552 (90)

- superpositions - how do we characterize phase?



- interferometer



## 2-qubit quantum logic – overview:

- 2-qubit gates  $\equiv$  conditional dynamics
  - e.g. “flip qubit #2 iff. qubit #1 =  $|1\rangle$ ”
- need coupling between the qubits
 

(also need coupling for entanglement...)

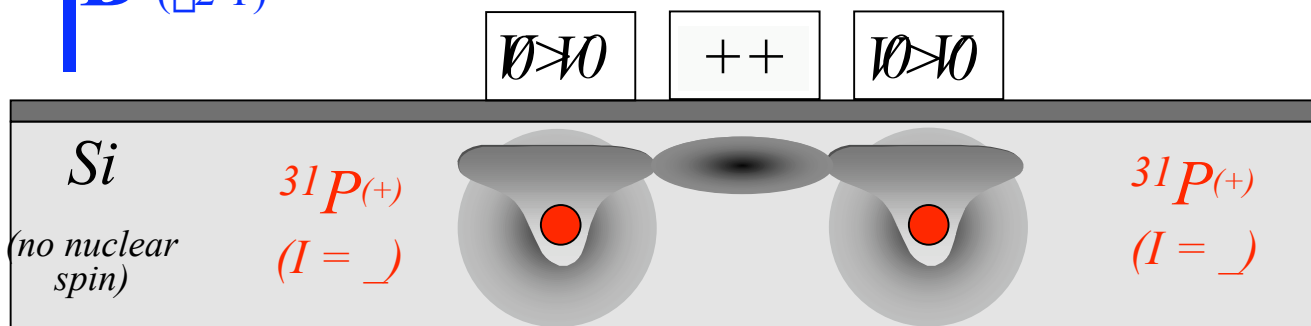
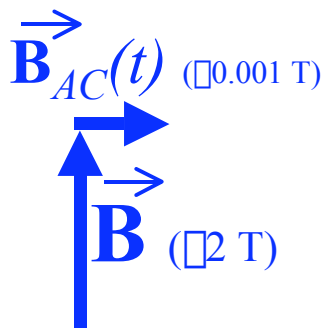
e.g.

  - Coulomb interaction
  - dipole-dipole coupling
  - exchange interaction, Fermi exchange...
- must be able to (effectively) turn coupling ON/OFF!

# Solid-state qubits (Kane)

B.E. Kane, Nature 393, 133 (98)

- nuclear spin qubits:  $\mathbf{B}$  splits the energy of  $| \downarrow \downarrow \rangle, | \uparrow \downarrow \rangle$ 
  - energies modified by  $e^-$  distributions (hyperfine)
  - use chip technology to change  $e^-$  distributions

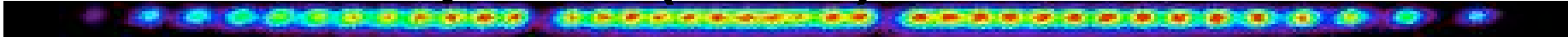


$0 < V < E_{HF}$   
 $\mathbf{B}_{AC}$  not resonant

$V > 0 > E_{HF}'$   
 $\mathbf{B}_{AC}$  resonant

$U > 0$   
 $e^- e^-$  interaction  
 $\square$  nucl. interac.

# Solid-state qubits (Kane)



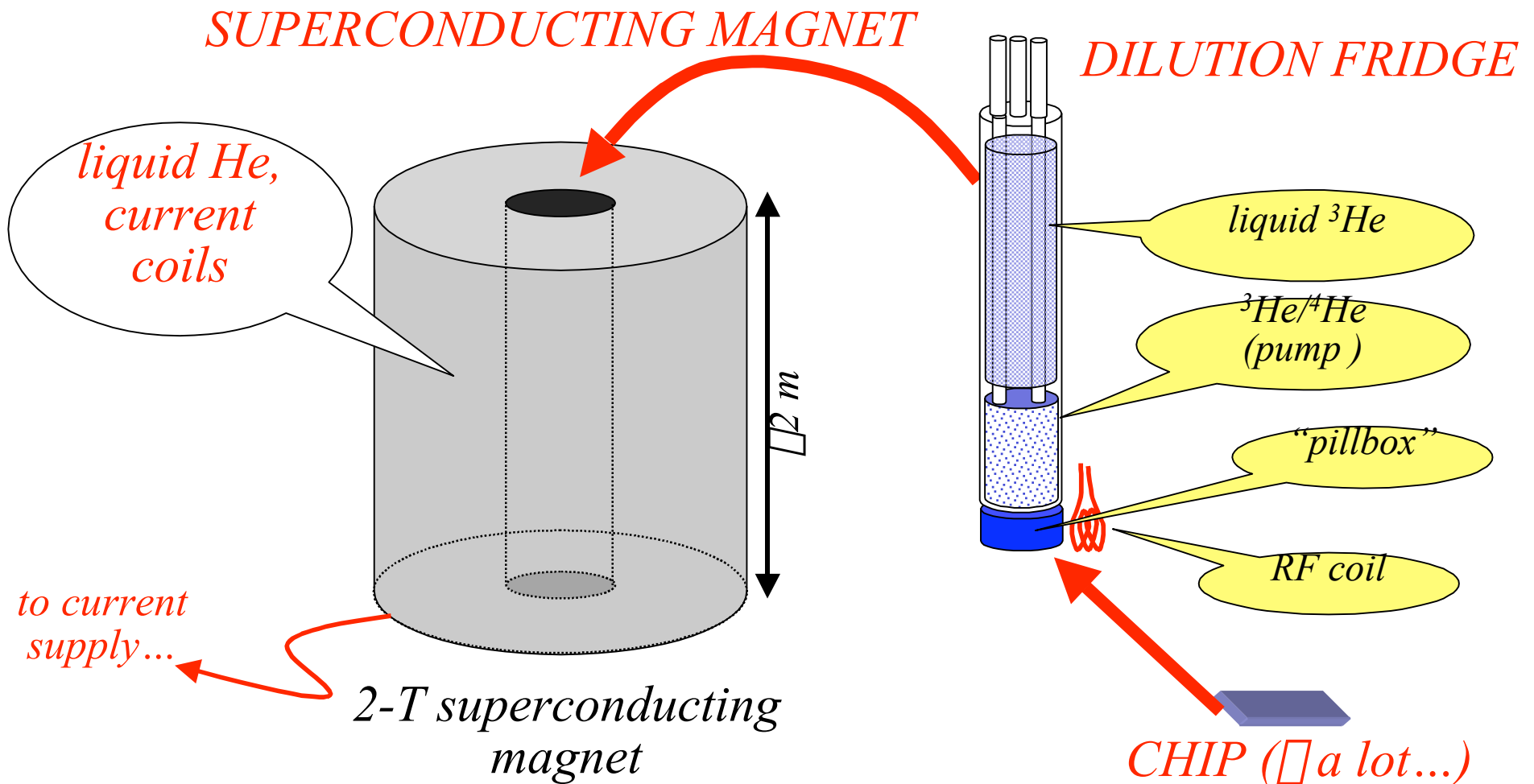
- initialization:
  - run at  $k_B T \ll \hbar \nu_B \hbar \quad T < 100 \text{ mK}, B > 2 T$
- readout:
  - map nuclear spin onto  $e^\uparrow$  spin
  - apply gate voltage
  - only get current if  $e^\uparrow$  are in singlet (  $\uparrow \downarrow \downarrow \uparrow$  )
- problems?
  - donor placement
  - decoherence mechanisms in bulk? (charge noise)
  - readout?



# Kane, cont'd.:

- So what will it look like?

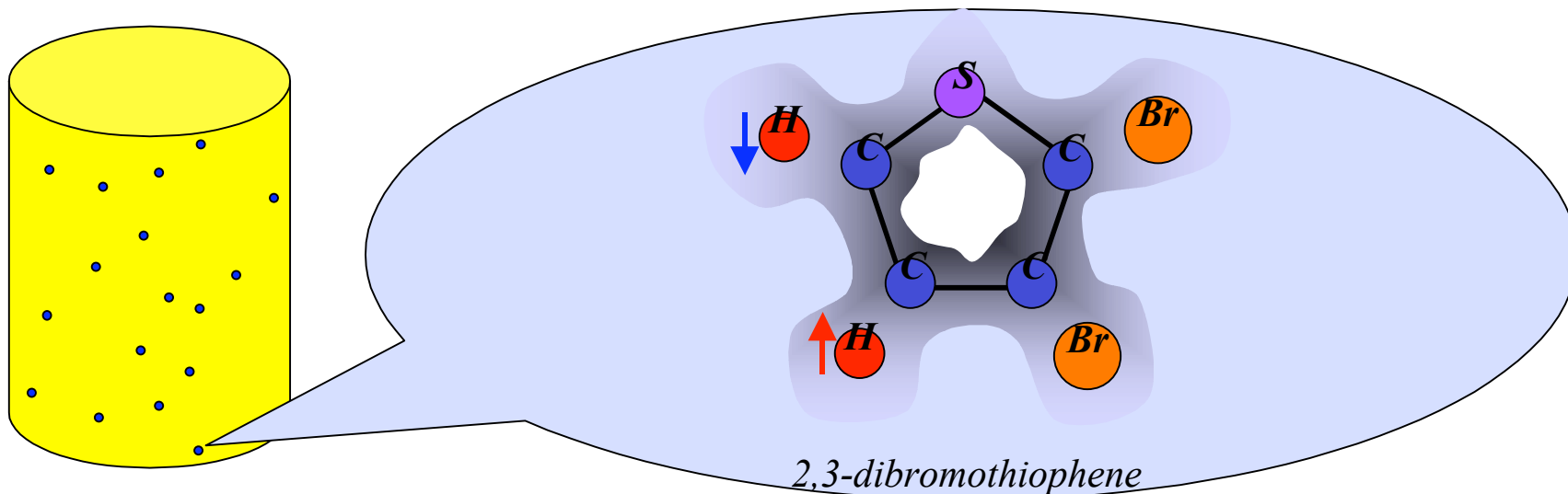
\*...I have *no* idea... (“dammit, Jim, I’m an atomic physicist, not a condensed matter physicist!”)



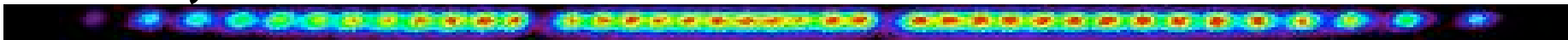
# Room-T, liquid-state NMR:

Cory, Fahmy, and Havel, Proc. Nat. Acad. Sci. USA **94**, 1634 (97); Gershenfeld and Chuang, Science **275**, 350 (97); Jones and Mosca, J. Chem. Phys. **109**, 1648 (98), Laflamme, Knill, Zurek, Catasti, and Mariappan, Phil. Trans. Roy. Soc. Lond. **A356**, 1941 (98).

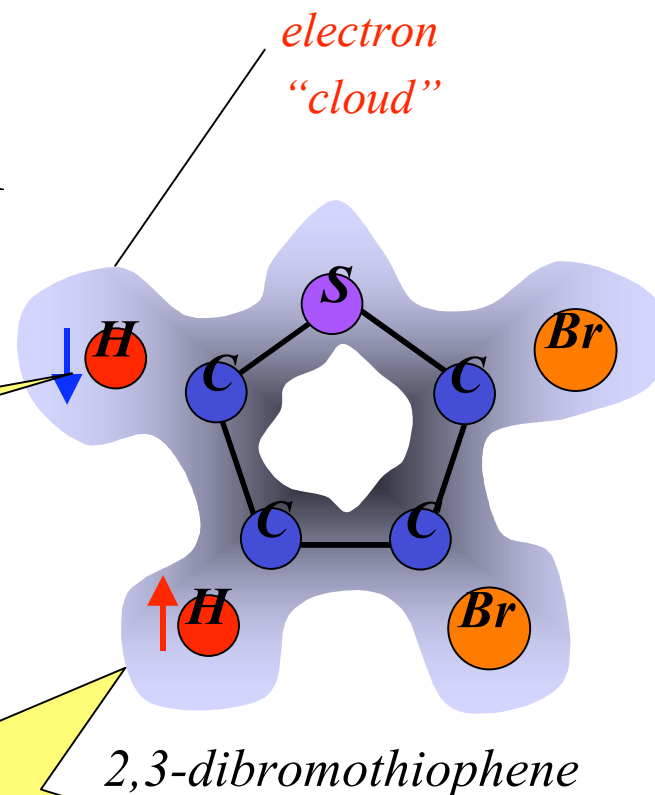
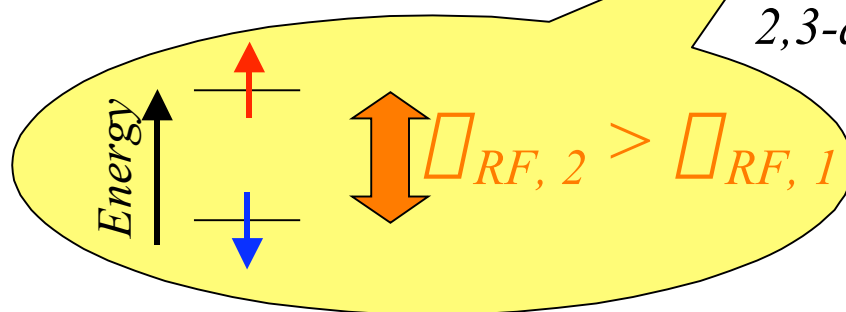
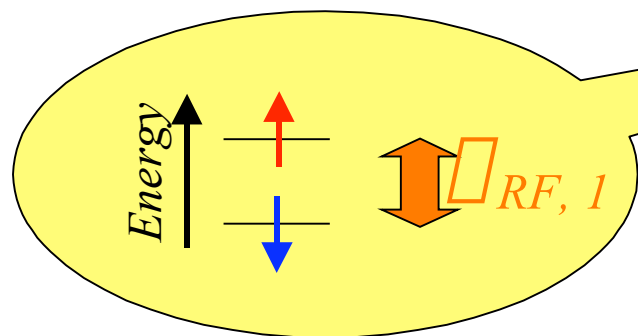
- ensemble of  $N \sim 10^{22}$  molecules
  - room temperature
  - liquid (□ molecular interactions negligible...)
- each molecule has  $n$  spin- $\frac{1}{2}$ 's (or higher)
  - e.g. hydrogen/protons, carbon, ...
- *each molecule is one quantum computer*



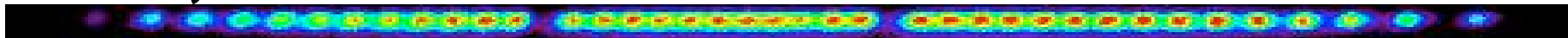
# NMR, cont'd.:



- apply **B** to split *H* spin energy levels
- apply **B<sub>RF</sub>** to change *H* spin
- *H* energies depend on location
  - due to coupling to other atoms



# NMR, cont'd.:

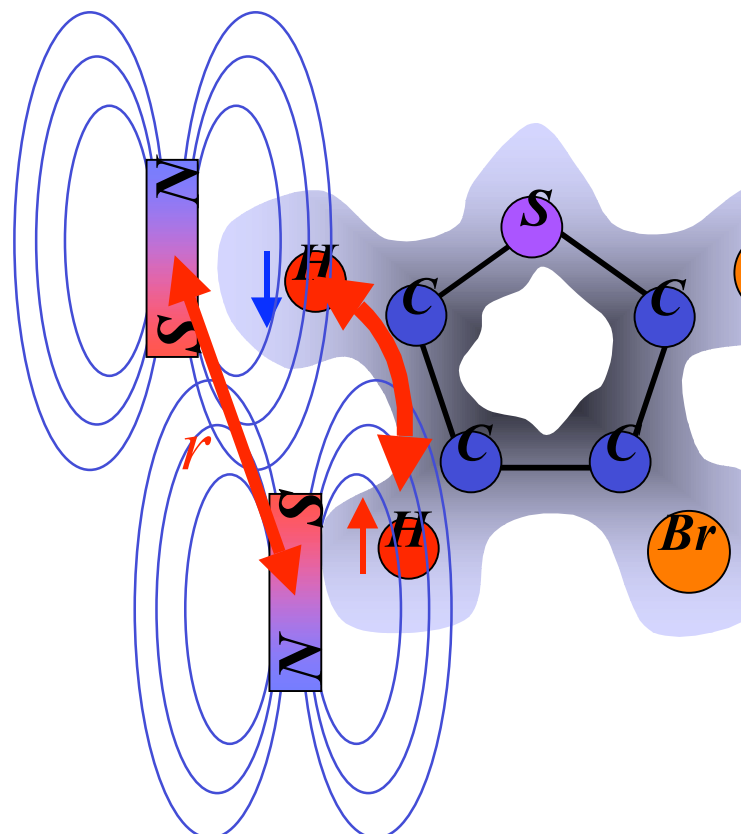


- Coupling:

$$H_{int} = \frac{-\mu_1 \mu_2}{4 r^3} [\mu_1 \cdot \mu_2 - 3(\mu_1 \cdot n)(\mu_2 \cdot n)] + \frac{J}{8} [\mu_1 \cdot \mu_2]$$

## Dipolar coupling

- averages to zero if molecules are tumbling



## J-coupling

- molecular orbitals:
- atoms share electrons through chemical bonds

$$\mu J \mu_{z1} \mu_{z2} / 4$$

# NMR Quantum logic:

- single-qubit gates:
  - apply  $\mathbf{B}_{RF}$
- two-qubit gates:
  - *don't* do nothing!...
  - J-coupling is always on
  - but

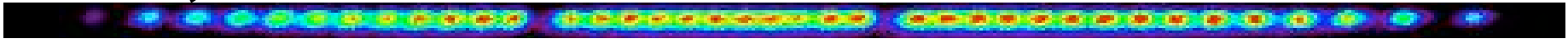
$$e^{-i2\pi\nu_x t} e^{-i\alpha\nu_z t} e^{-i2\pi\nu_x t} = e^{+i\alpha\nu_z t}$$

$$\square e^{-iJ\nu_z \nu_z t/4} e^{-i2\pi\nu_x t} e^{-iJ\nu_z \nu_z t/4} e^{-i2\pi\nu_x t} = e^{-i2\pi\nu_z t}$$

*“refocussing”*

*- turn off for quantum logic!*

# NMR, cont'd.:



- “but NMR is done at room T!?!?!...”

## 1. “effective pure state”

- thermodynamics  $\square$  less population in higher-energy states

$$\square \square \square \square \uparrow + \square U |00\dots 0\rangle \square \square |00\dots 0\rangle U^\dagger$$

(cost in resources...)

- *readout insensitive to diagonal terms in  $\square$*

$\square$  *only read out  $|00\dots 0\rangle \square \square |00\dots 0\rangle$  part*

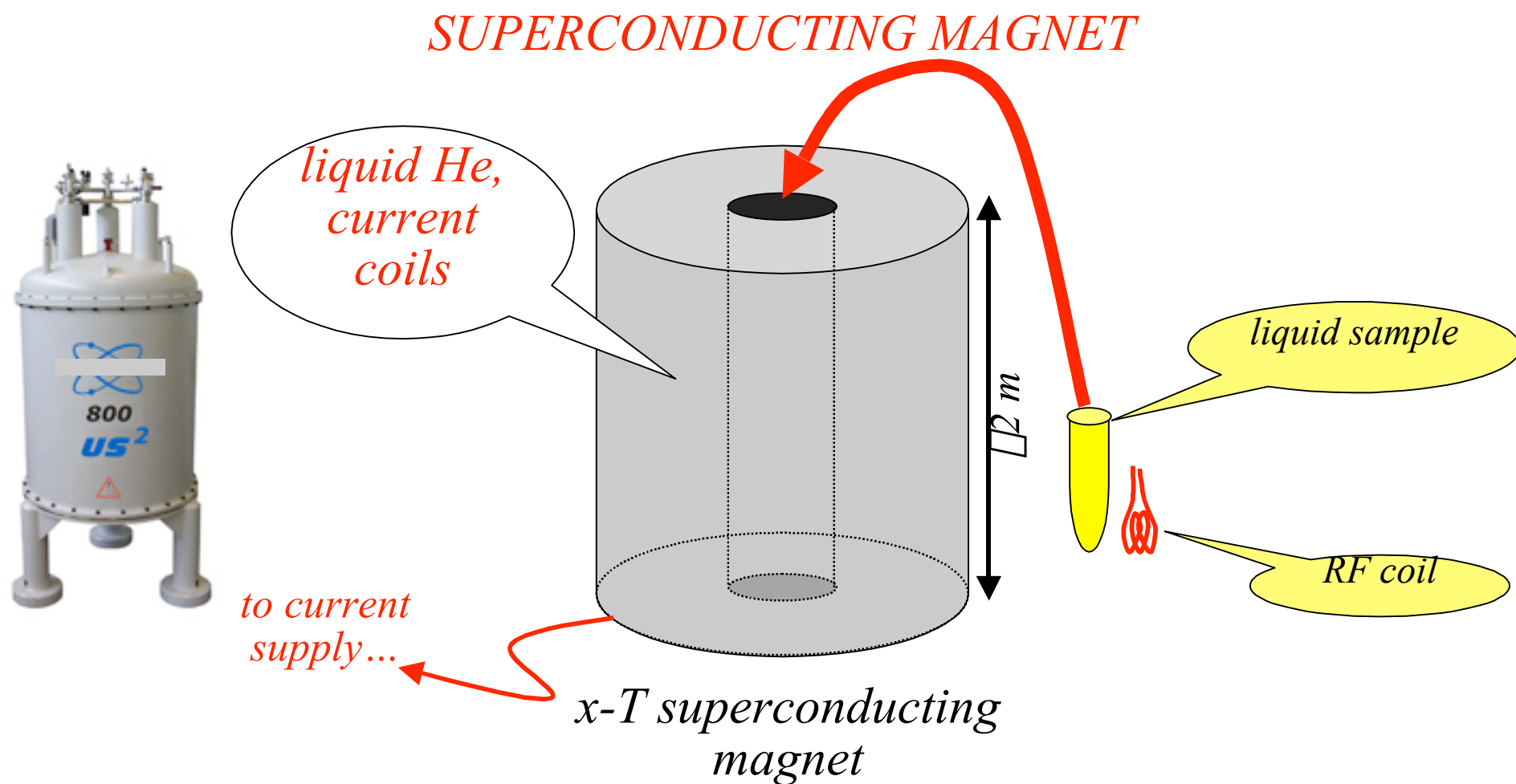
## 2. ensemble readout

- ensemble averages, but still works for some probs.

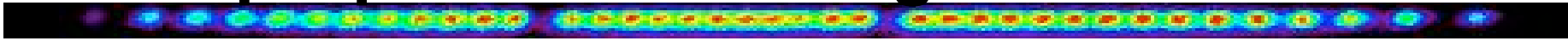
# NMR, cont'd.:

- So what does it look like?

\*...I have *no* idea... (dammit, Jim, I'm an atomic physicist, not a condensed matter physicist!)



# Other proposed technologies:

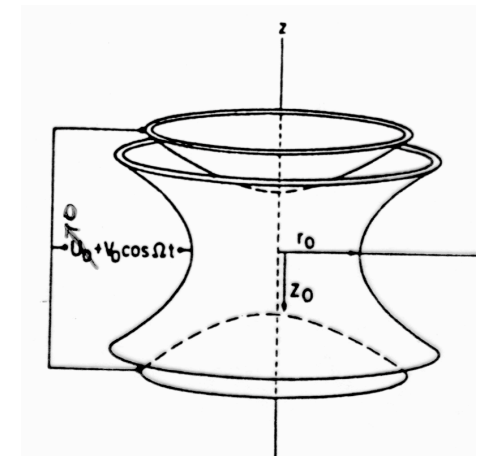
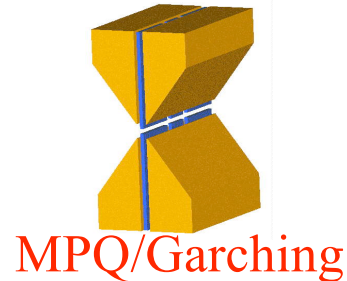
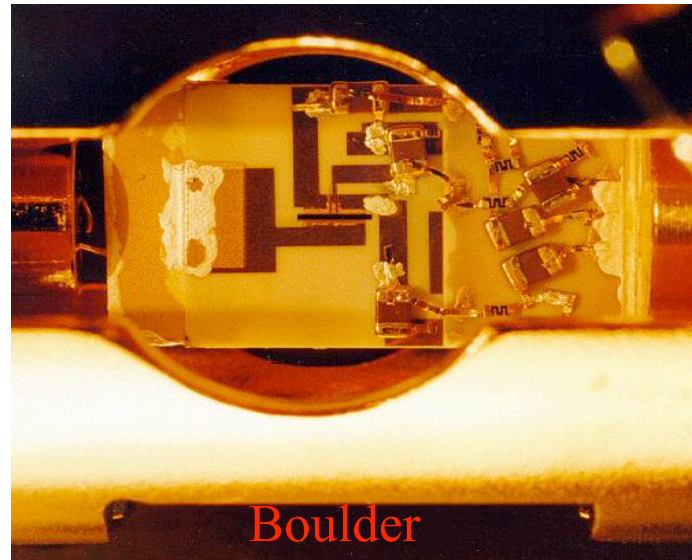
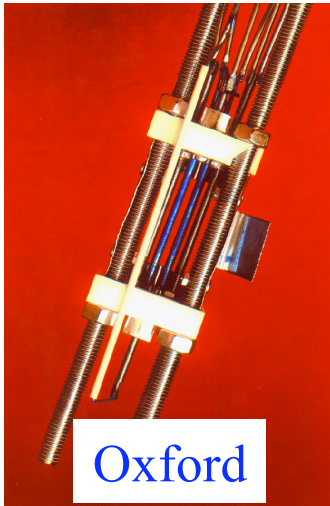


- recall
  - strong, switchable, controllable qubit interactions
  - no other interactions!
    - *photons* (see also *R. L.Fl.!*)
    - ?← • Josephson junctions (Martinis, et al., Phys. Rev. Lett. 89, 117901 (02) and references therein...[Nakamura, Devoret, Martinis, Han,...])
    - ?← • electron dots (Bayer *et al.*, Science 291, 451 (03) and references therein...[Steel, Sherwin, Hawrylak,...])
      - electrons floating above liquid He
      - etc...
    - **trapped atomic ions**



# Ion traps for quantum computing:

- store quantum information inside atoms
- build on existing quality of superpositions
- need way to implement 2-qubit gates
- plus a lot more!...



G. Werth, *Progress in Atomic Spectroscopy*,  
H.J. Beyer, H. Kleinpoppen, eds