

## Routing of Optical States by Atomic Media

Jürgen Appel, Frank Vewinger, Eden Figueroa, Georg Günter,  
Peter Marzlin und Alexander I. Lvovsky

Institute for Quantum Information Science  
University of Calgary, Canada

18. 8. 2006

Canadian Quantum Information Students' Conference



Institute for  
**Quantum Information Science**  
at the University of Calgary

# Introduction

## Routing of Optical States by Atomic Media

Problem:



- Non-linear upconversion in crystals (Tanzilli et al. Nature **437** (2005))  
→ low conversion efficiency
- Storage in atomic vapor, retrieval on second transition (Zibrov et al. PRL **88** (2002))  
→ Losses due to storage
- Four wave mixing  
→ created fields are absorbed

# Introduction

## Routing of Optical States by Atomic Media

Problem:



- Non-linear upconversion in crystals (Tanzilli et al. Nature **437** (2005))  
→ low conversion efficiency
- Storage in atomic vapor, retrieval on second transition (Zibrov et al. PRL **88** (2002))  
→ Losses due to storage
- Four wave mixing  
→ created fields are absorbed

Our Approach:

- Transfer the optical fields adiabatically (similar to STIRAP)
- Avoid absorption using **electromagnetically induced transparency**

# Introduction

## Routing of Optical States by Atomic Media

Problem:

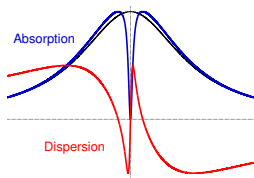
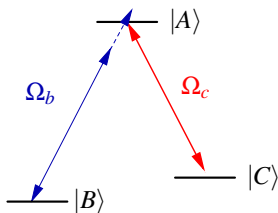


- Non-linear upconversion in crystals (Tanzilli et al. Nature **437** (2005))  
→ low conversion efficiency
- Storage in atomic vapor, retrieval on second transition (Zibrov et al. PRL **88** (2002))  
→ Losses due to storage
- Four wave mixing  
→ created fields are absorbed

Our Approach:

- Transfer the optical fields adiabatically (similar to STIRAP)
- Avoid absorption using **electromagnetically induced transparency**
- Framework for multiple input modes

# Electromagnetically Induced Transparency: Slow Light



EIT is a quantum interference effect.  
Atomic medium with  $\Lambda$ -like level configuration

$\Omega_c$  control field

$\Omega_b$  (weak) signal field

$$\hat{H}_{\text{int}} = \hbar \Omega_b |A\rangle\langle B| + \hbar \Omega_c |A\rangle\langle C| + \text{h.a.}$$

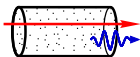
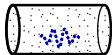
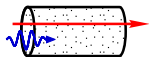
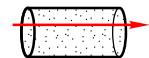
$$|\text{Dark state}\rangle \propto \Omega_c |B\rangle - \Omega_b |C\rangle$$

- Excitation processes  $|C\rangle \rightarrow |A\rangle$  and  $|B\rangle \rightarrow |A\rangle$  interfere destructively.
- No absorption on two-photon resonance.
- Sensitive phase dependence.
- Strong dispersion  $\frac{dn}{d\omega}$ :

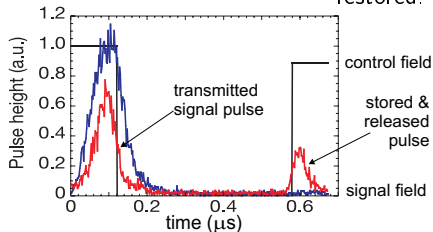
$$\Rightarrow \text{Reduction of group velocity } v_g = \frac{c}{n + \omega \frac{dn}{d\omega}}$$

# Electromagnetically Induced Transparency: Stopping of Light

## Stopping of light



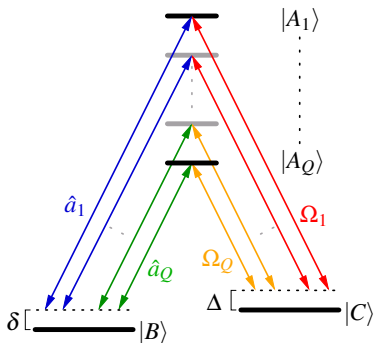
- Turn on control laser: medium gets transparent for signal mode
- Signal enters the medium: pulse is slowed down, compressed spatially
- Control is turned off adiabatically: group velocity reduces to zero, pulse is stored in a collective ground state superposition.
- Control is turned on adiabatically: signal pulse is restored.



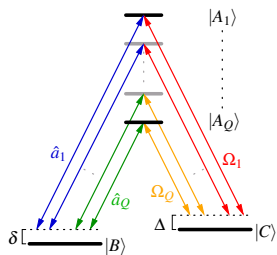
Reproduced from T. Chanelière et. al., Nature 438 (2006)

EIT in Multi- $\Lambda$  Systems

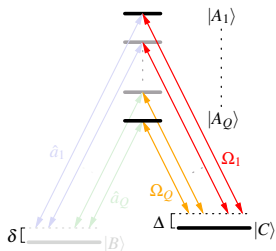
PRA 73, 013804 (2006)

 $\Lambda$ -system with multiple excited levels

- $\Omega_q$  strong  $\Rightarrow$  Classical control fields
  - $\hat{a}_q$  weak quantized signal field modes
  - $\delta, \Delta$  small detunings  
 $\Rightarrow$  each field couples one transition
  - Dicke limit: signals' coupling strengths  $g_q$  identical for all atoms
- How to handle EIT in this system?
  - Which signal mode will propagate?

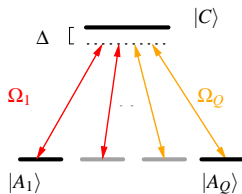
Multi- $\Lambda$  system



Multi- $\Lambda$  system

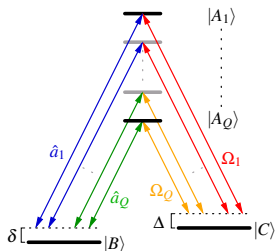
Interaction Hamiltonian  
(neglecting weak fields  $\hat{a}_q$ ):

$$\hat{H}_{\text{int}}^{(0)} = -\hbar\Delta - \hbar \sum_q \Omega_q |A_q\rangle \langle C| + \text{H.a.}$$



$Q - 1$  excited 'dark states'  $|ED_r\rangle$

one excited 'bright state'  $|EB\rangle$

Multi- $\Lambda$  system

Only one superposition of the excited states

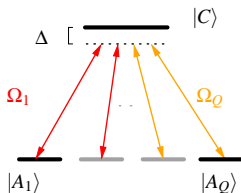
$$|EB\rangle = \sum \frac{\Omega_q}{\Omega} |A_q\rangle$$

couples to  $|C\rangle$ .

$$\Omega = \sqrt{\sum |\Omega_q|^2}$$

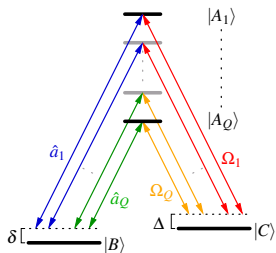
Interaction Hamiltonian  
(neglecting weak fields  $\hat{a}_q$ ):

$$\hat{H}_{\text{int}}^{(0)} = -\hbar\Delta - \hbar \sum_q \Omega_q |A_q\rangle \langle C| + \text{H.a.}$$



$Q - 1$  excited 'dark states'  $|ED_r\rangle$

one excited 'bright state'  $|EB\rangle$

Multi- $\Lambda$  system

Only one superposition of the excited states

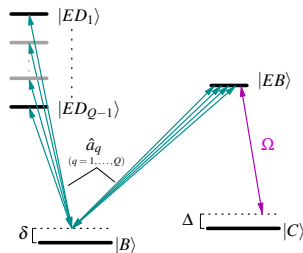
$$|EB\rangle = \sum \frac{\Omega_q}{\Omega} |A_q\rangle$$

couples to  $|C\rangle$ .

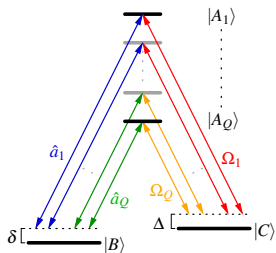
$$\Omega = \sqrt{\sum |\Omega_q|^2}$$

## Change of the atomic basis:

$$\{|A_q\rangle\} \rightarrow \{|EB\rangle, |ED_r\rangle\}$$



Quantum fields  $\hat{a}_q$  couple to all  $|ED_r\rangle \rightarrow$  Absorption

Multi- $\Lambda$  system

Only one superposition of the excited states

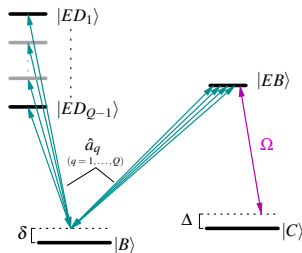
$$|EB\rangle = \sum \frac{\Omega_q}{\Omega} |A_q\rangle$$

couples to  $|C\rangle$ .

$$\Omega = \sqrt{\sum |\Omega_q|^2}$$

## Change of the atomic basis:

$$\{|A_q\rangle\} \rightarrow \{|EB\rangle, |ED_r\rangle\}$$

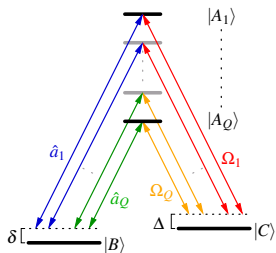


Quantum fields  $\hat{a}_q$  couple to all  $|ED_r\rangle \rightarrow$  Absorption

Only one superposition couples only to  $|EB\rangle$ :

$$\hat{b}_Q = \frac{1}{R} \sum \frac{\Omega_q^*}{g_q} \hat{a}_q.$$

$$R = \sqrt{\sum |\Omega_q/g_q|^2}$$

Multi- $\Lambda$  system

Only one superposition of the excited states

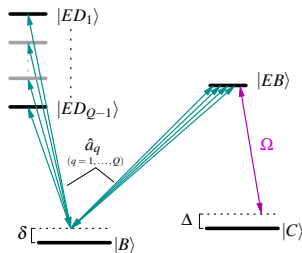
$$|EB\rangle = \sum \frac{\Omega_q}{\Omega} |A_q\rangle$$

couples to  $|C\rangle$ .

$$\Omega = \sqrt{\sum |\Omega_q|^2}$$

## Change of the atomic basis:

$$\{|A_q\rangle\} \rightarrow \{|EB\rangle, |ED_r\rangle\}$$



Quantum fields  $\hat{a}_q$  couple to all  $|ED_r\rangle \rightarrow$  Absorption

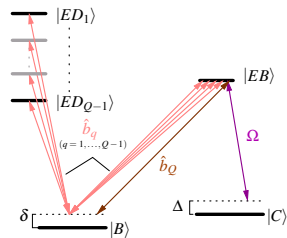
Only one superposition couples only to  $|EB\rangle$ :

$$\hat{b}_Q = \frac{1}{R} \sum \frac{\Omega_q^*}{g_q^*} \hat{a}_q.$$

$$R = \sqrt{\sum |\Omega_q/g_q|^2}$$

## Change of the optical basis:

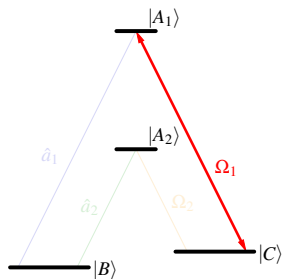
$$\{\hat{a}_q\} \rightarrow \{\hat{b}_q\}$$



Reduction to single- $\Lambda$  system:

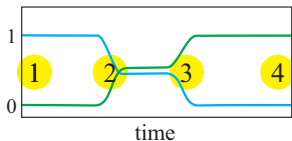
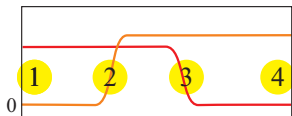
- One mode  $\hat{b}_Q$  experiences EIT
- The other modes are absorbed

# Raman Adiabatic Transfer of Optical States

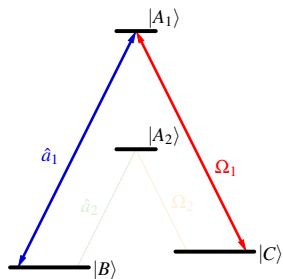


Consider performing the following protocol in a medium with two excited levels.

- Switch on control field  $\Omega_1$   
EIT-mode  $\hat{b}_2(\Omega_1, \Omega_2) = \hat{a}_1$ .

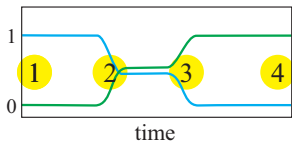
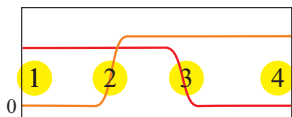


# Raman Adiabatic Transfer of Optical States

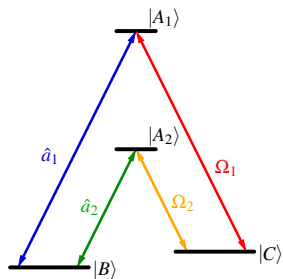


Consider performing the following protocol in a medium with two excited levels.

- Switch on control field  $\Omega_1$   
EIT-mode  $\hat{b}_2(\Omega_1, \Omega_2) = \hat{a}_1$ .  
Couple a signal into mode  $\hat{a}_1$ .



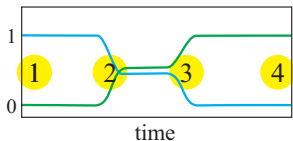
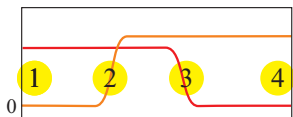
# Raman Adiabatic Transfer of Optical States



Consider performing the following protocol in a medium with two excited levels.

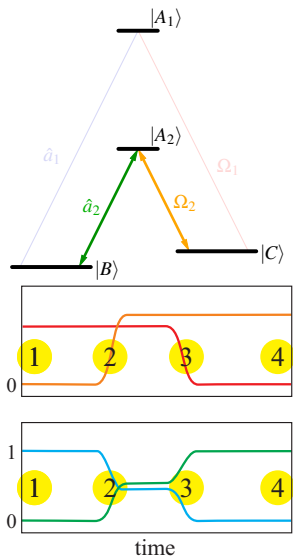
- 1 Switch on control field  $\Omega_1$   
EIT-mode  $\hat{b}_2(\Omega_1, \Omega_2) = \hat{a}_1$ .  
Couple a signal into mode  $\hat{a}_1$ .
- 2 Slowly switch on control field  $\Omega_2$   
signal follows into the new EIT-Mode adiabatically

$$\hat{b}_2(\Omega_1, \Omega_2) \rightarrow \frac{1}{R} \left( \frac{\Omega_1^*}{g_1^*} \hat{a}_1 + \frac{\Omega_2^*}{g_2^*} \hat{a}_2 \right)$$





# Raman Adiabatic Transfer of Optical States



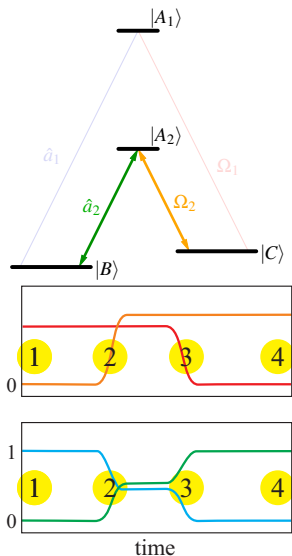
Consider performing the following protocol in a medium with two excited levels.

- 1 Switch on control field  $\Omega_1$   
EIT-mode  $\hat{b}_2(\Omega_1, \Omega_2) = \hat{a}_1$ .  
Couple a signal into mode  $\hat{a}_1$ .
- 2 Slowly switch on control field  $\Omega_2$   
signal follows into the new EIT-Mode adiabatically

$$\hat{b}_2(\Omega_1, \Omega_2) \rightarrow \frac{1}{R} \left( \frac{\Omega_1^*}{g_1^*} \hat{a}_1 + \frac{\Omega_2^*}{g_2^*} \hat{a}_2 \right)$$

- 3 Switch off control field  $\Omega_1$ :  
EIT-mode  $\hat{b}_2(\Omega_1, \Omega_2) \rightarrow \hat{a}_2$ .

# Raman Adiabatic Transfer of Optical States

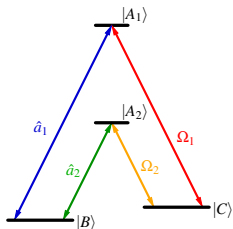


Consider performing the following protocol in a medium with two excited levels.

- 1 Switch on control field  $\Omega_1$   
EIT-mode  $\hat{b}_2(\Omega_1, \Omega_2) = \hat{a}_1$ .  
Couple a signal into mode  $\hat{a}_1$ .
- 2 Slowly switch on control field  $\Omega_2$   
signal follows into the new EIT-Mode adiabatically

$$\hat{b}_2(\Omega_1, \Omega_2) \rightarrow \frac{1}{R} \left( \frac{\Omega_1^*}{g_1^*} \hat{a}_1 + \frac{\Omega_2^*}{g_2^*} \hat{a}_2 \right)$$

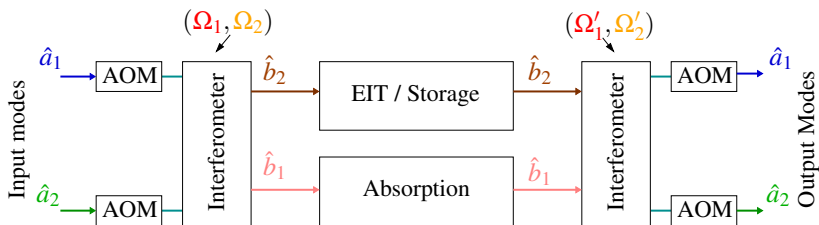
- 3 Switch off control field  $\Omega_1$ :  
EIT-mode  $\hat{b}_2(\Omega_1, \Omega_2) \rightarrow \hat{a}_2$ .
- 4 The original signal has been converted from mode  $\hat{a}_1$  to mode  $\hat{a}_2 \Rightarrow$  STIRAP for optical modes



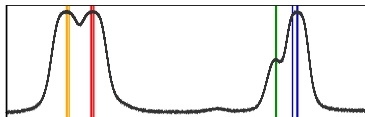
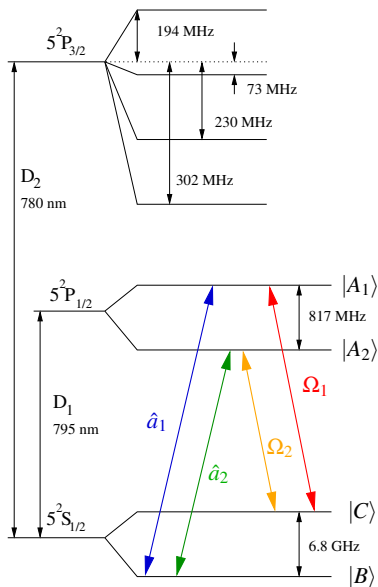
is equivalent to

## Adiabatic control allows routing

- One superposition of input modes is selected:  $\hat{b}_2(\Omega_1, \Omega_2)$ .
- This mode experiences EIT; orthogonal modes get absorbed.
- Phase and amplitudes of the input- and output modes are controlled by the control fields  $\Omega_1, \Omega_2$ .
- Equivalent to a linear optics circuit

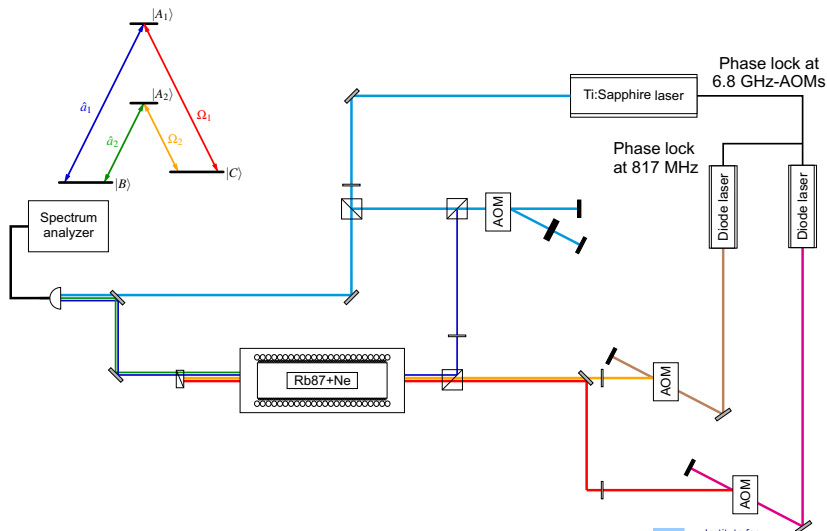


# The Physical System

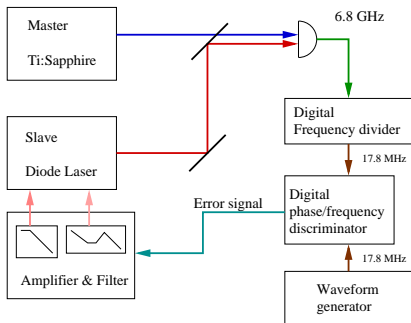


- Atomic Medium:  $^{87}\text{Rb}$
- 50 mm long cell, 5 torr Neon as buffer gas, cell magnetically shielded,  $60^\circ\text{C}$
- $\hat{a}_1$  Ti:Sapphire laser
- $\hat{a}_2$  generated mode; detection via heterodyning to Ti:Sapphire
- $\Omega_1$  Diode laser, phase locked to Ti:Sapphire  $\hat{a}_1$  at 6.8 GHz
- $\Omega_2$  Diode laser, phase locked to  $\Omega_1$  at 817 MHz

# The Setup

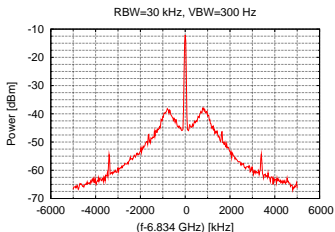


# Laser Phase Lock



## Properties

- using mobile communication ICs
- loop bandwidth  $>1$  MHz
- large capture range
- scan, sweep, modulation control
- very versatile



# The Experiment

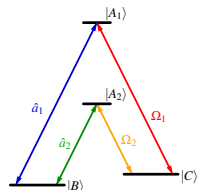
## Experimental procedure:

- Mode match all the beams for maximum beat signal
- Create four-wave-mixing: Shine in both control lasers and the signal

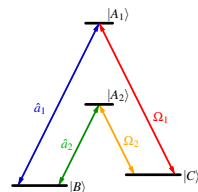
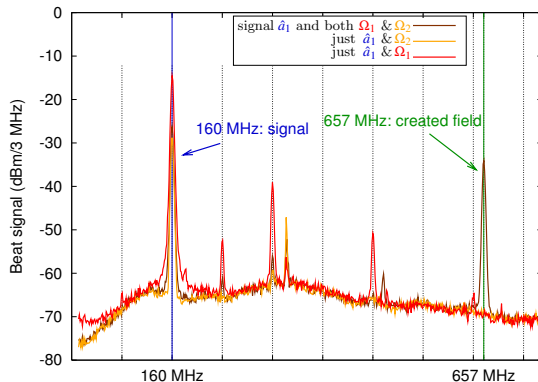
$$\rightarrow \omega_{\text{signal}_2} = \omega_{\text{signal}_1} - (\omega_{\text{pump}_1} - \omega_{\text{pump}_2})$$

is created.

- Store light pulses
- Retrieve stored pulses with second pump laser
- Observe RATOS



# Four-wave mixing



160 MHz

Beat note between  
signal and Ti:Sa

657 MHz

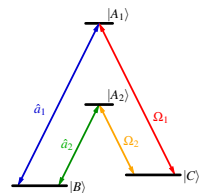
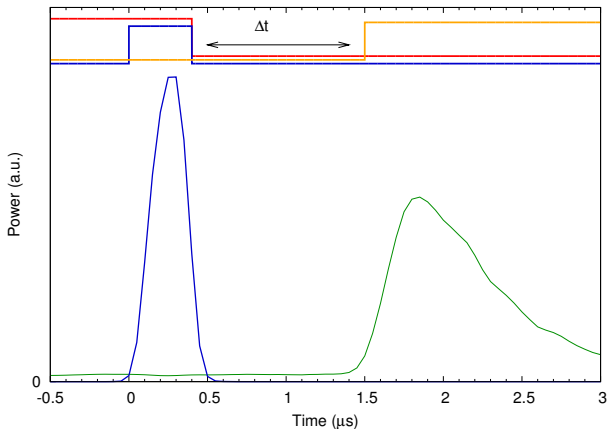
Beat note between  
created field and Ti:Sa

⇒ Mode matching is sufficient for the creation and detection of the new field



# Storage of light

Storage and retrieval of a 400 ns pulse.

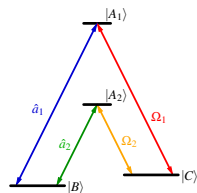
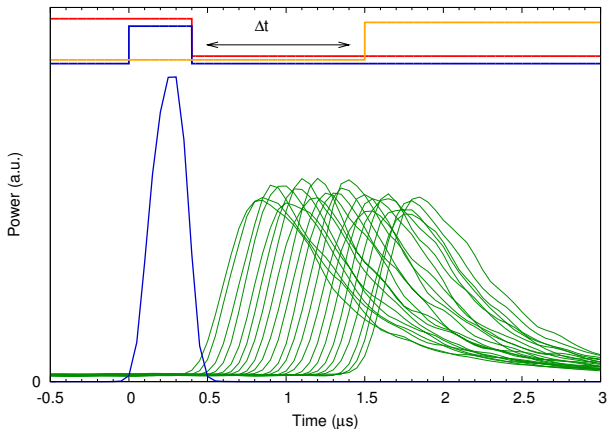


- Creation of ground state coherence
- Transfer efficiency  $> 80\%$  compared to retrieval on same frequency.

$\Rightarrow$  Frequency conversion is possible if the pulse is stored

# Storage of light

Storage and retrieval of a 400 ns pulse.

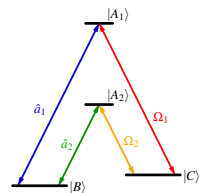
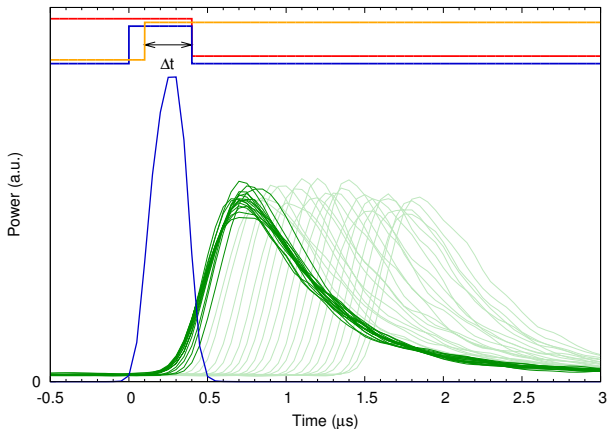


- Decoherence time  $\approx 100\mu\text{s}$ .
- Decoherence due to transit broadening

$\Rightarrow$  Frequency conversion is possible if the pulse is stored

# Adiabatic transfer

Adiabatic transfer of a 400 ns pulse.

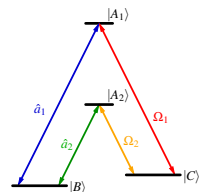
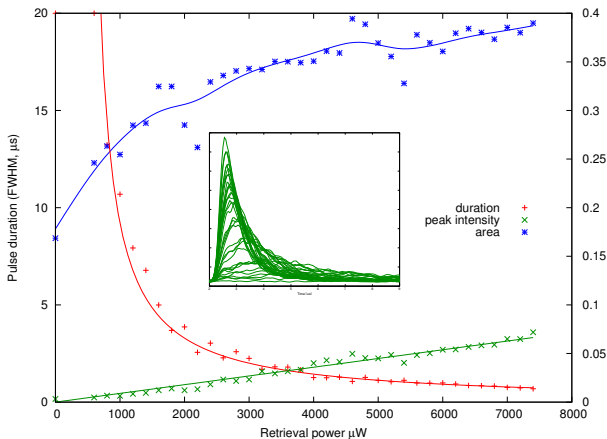


Adiabatic process

- Efficiency independent of overlap

# Adiabatic transfer

Adiabatic transfer of a 400 ns pulse.

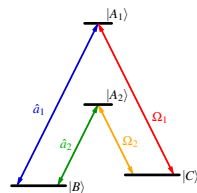
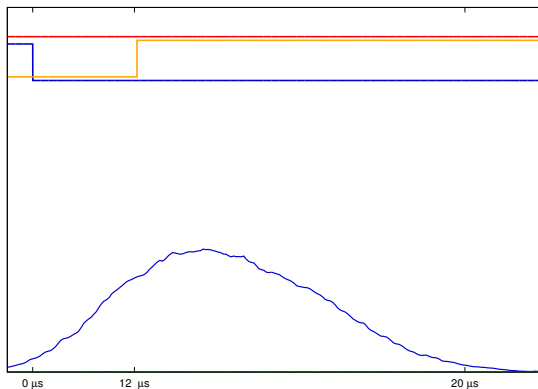


Adiabatic process

- Efficiency independent of overlap
- Efficiency independent of power

# Beam Splitting

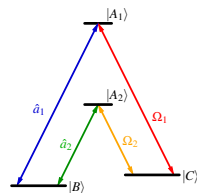
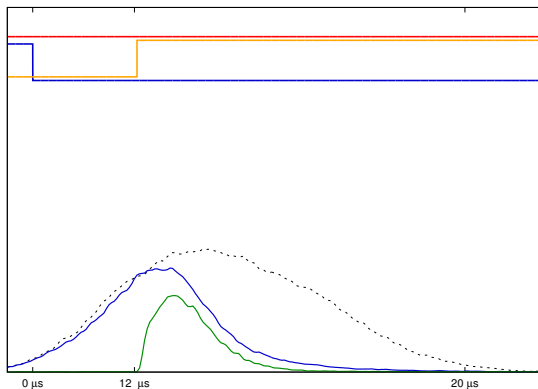
Control Powers:  $\Omega_1 \hat{=} 2.8 \text{ mW}$ ,  $\Omega_2 \hat{=} 0 \text{ mW}$



- Simultaneous retrieval with both control lasers splits the pulse into different frequency modes
- splitting ratio defined by  $\frac{\Omega_1}{\Omega_2}$

# Beam Splitting

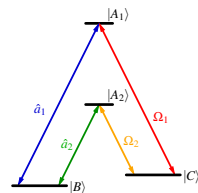
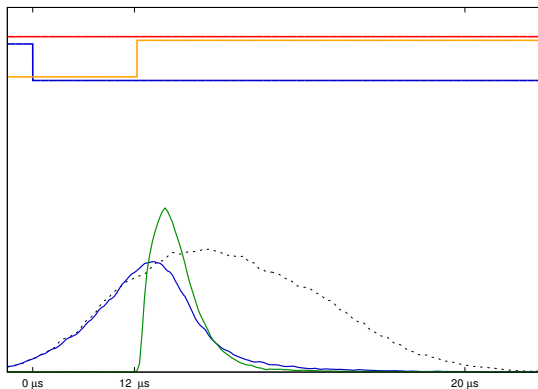
Control Powers:  $\Omega_1 \hat{=} 2.8 \text{ mW}, \Omega_2 \hat{=} 2.8 \text{ mW}$



- Simultaneous retrieval with both control lasers splits the pulse into different frequency modes
- splitting ratio defined by  $\frac{\Omega_1}{\Omega_2}$

# Beam Splitting

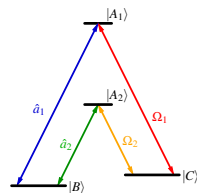
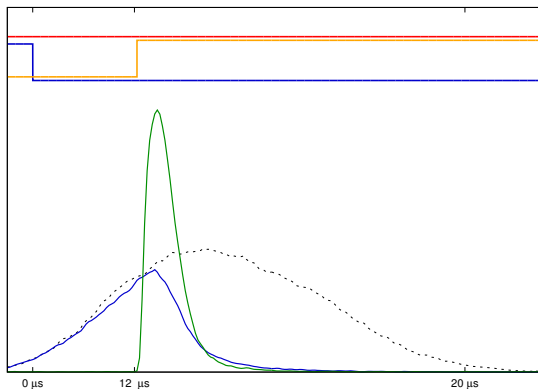
Control Powers:  $\Omega_1 \hat{=} 2.8 \text{ mW}, \Omega_2 \hat{=} 3.9 \text{ mW}$



- Simultaneous retrieval with both control lasers splits the pulse into different frequency modes
- splitting ratio defined by  $\frac{\Omega_1}{\Omega_2}$

# Beam Splitting

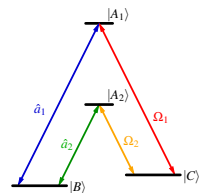
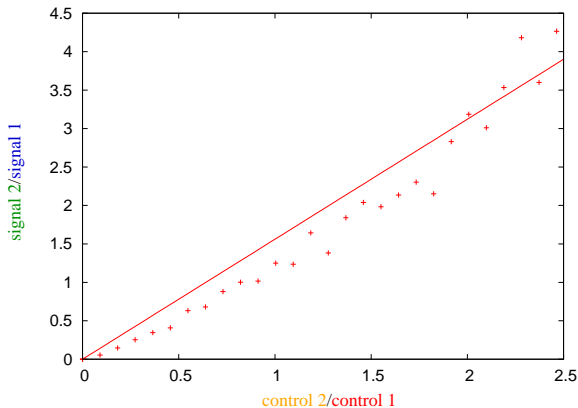
Control Powers:  $\Omega_1 \hat{=} 2.8 \text{ mW}, \Omega_2 \hat{=} 6.5 \text{ mW}$



- Simultaneous retrieval with both control lasers splits the pulse into different frequency modes
- splitting ratio defined by  $\frac{\Omega_1}{\Omega_2}$

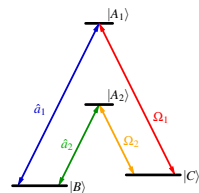
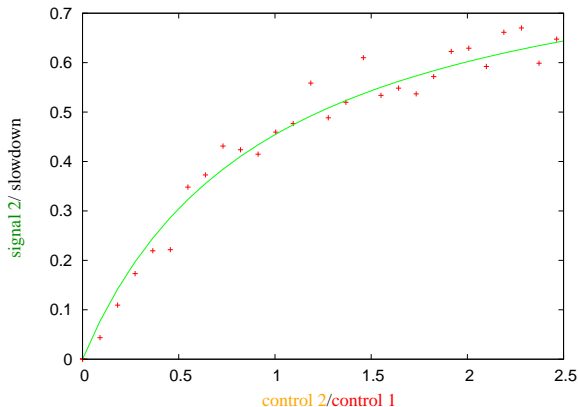


# Beam Splitting



- Simultaneous retrieval with both control lasers splits the pulse into different frequency modes
- splitting ratio defined by  $\frac{\Omega_1}{\Omega_2}$

# Beam Splitting

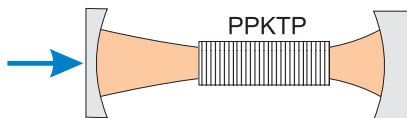


- Simultaneous retrieval with both control lasers splits the pulse into different frequency modes
- splitting ratio defined by  $\frac{\Omega_1}{\Omega_2}$

# Conclusion

- **Raman Adiabatic Transfer of Optical States**
  - Theory for an arbitrary number of upper levels  
⇒ Multiport beamsplitter possible
  - Experimental realization in  $^{87}\text{Rb}$ , D1 line  
⇒ 2 upper levels
  - Experimentally robust; only two-photon resonance must be maintained
  - EIT condition suppresses absorption of the created field
  - Applications
    - All-optical routing
    - Frequency conversion
    - Quantum state engineering
- To Do:
  - Show that RATOS transfers the quantum state
  - Transfer squeezed states
  - Transfer Fock states
  - Couple in a two-mode state

# Source of non-classical light for experiments with atoms



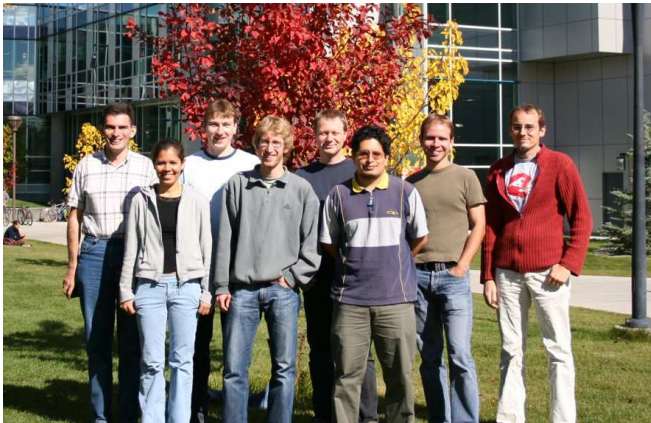
## Optical Parametric Amplifier

- 5 mm long PPKTP crystal
- type-I configuration
- pumped by frequency doubled Ti:Sa
- 75 mm linear resonator
- doubly resonant
- oscillation threshold  $\approx 50$  mW

## Properties

- narrow bandwidth ( $\approx 10$  MHz)
- high brightness
- resonant to atomic transition
- squeezed light close to threshold
- heralded Fock states far below threshold

# Thanks



Ph.D. & Postdoc positions available  
<http://qis.ucalgary.ca/quantech/>

Frank Vewinger,  
Eden Figueroa,  
Georg Günter,  
Peter Marzlin,  
Alex Lvovsky

Funding:

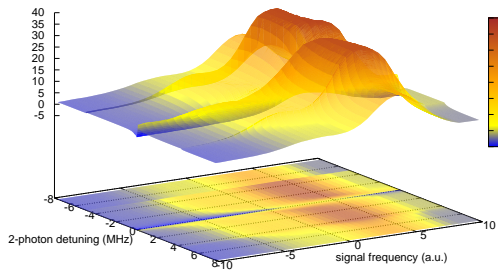
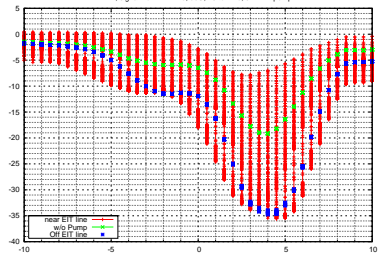
- CIAR
- NSERC
- AIF
- CFI

# TOC

- 1 Introduction
- 2 EIT
  - slow light & stopped light
  - multi- $\Lambda$  systems
  - mapping to single- $\Lambda$  system
- 3 RATOS
  - adiabatic process
  - linear optics equivalent
- 4 Implementation
  - level scheme
  - setup
  - phase lock
  - experiment
- 5 Results
  - 4-wave mixing
  - storage
  - RATOS
  - beam splitting
- 6 Conclusion

EIT,  $F=1$ EIT, signal from the  $F=1$  line, 55C, 10 torr Ne, 8.7 mW pump

Absorption (dB)

EIT, signal from  $F=1$  line, 55C, 10 torr Ne, 8.7 mW pump

EIT,  $F=2$ 